

Quantum frame covariance and subsystem relativity

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Okinawa Institute of Science and Technology



Quantum Information and Quantum Matter

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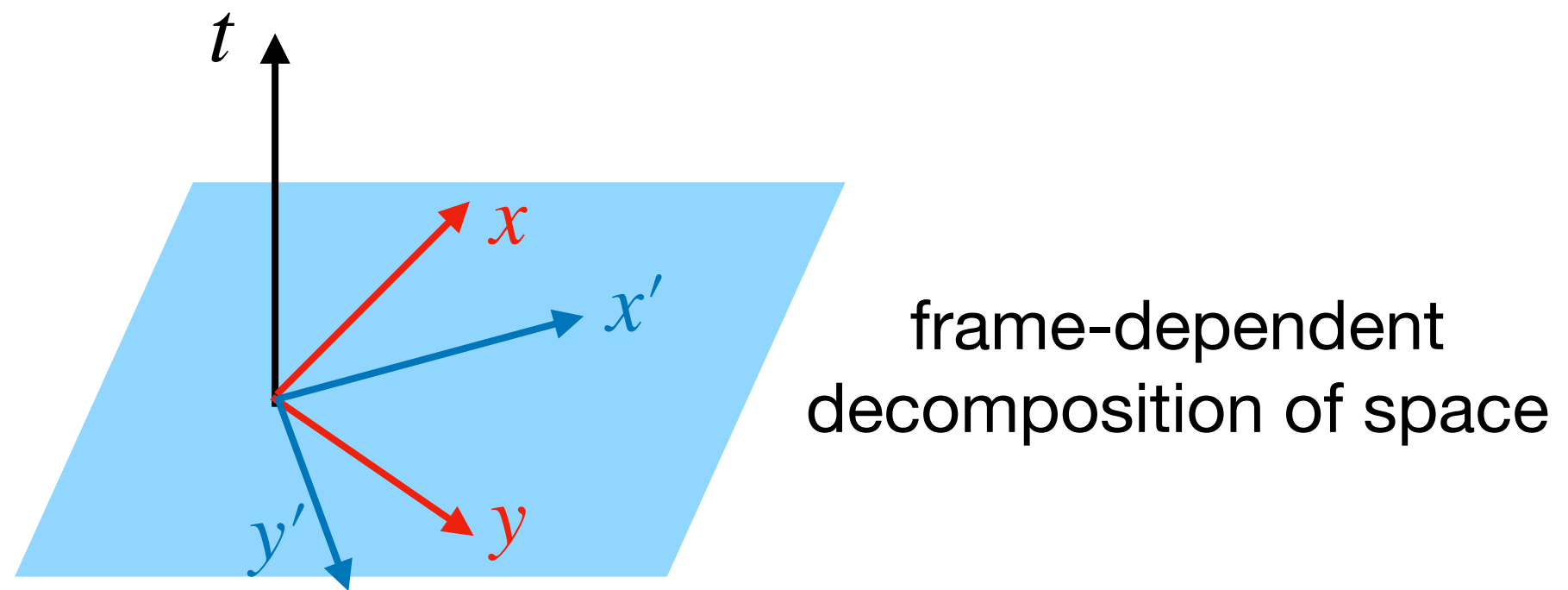
loosely based on:

Ahmad, Galley, PH, Lock, Smith, PRL 128 (2022) 170401
de la Hamette, Galley, PH, Loveridge, Müller 2110.13824
Vanrietvelde, PH, Giacomini, Castro-Ruiz, Quantum 4 (2020) 225
Mele, Kotecha, PH, *to appear soon*

Relativities

Galilean

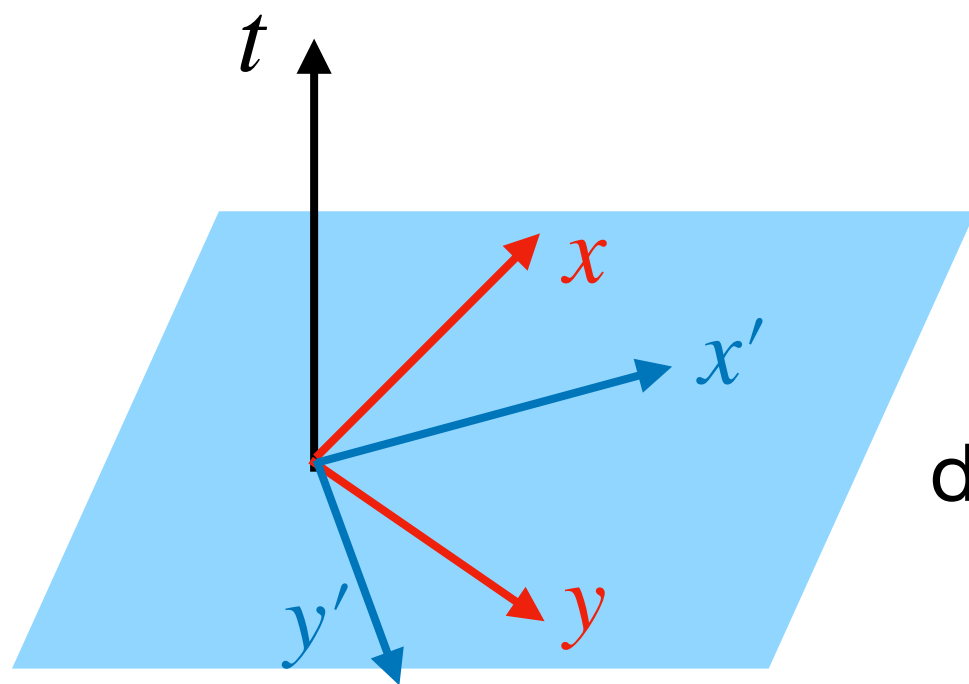
“All the laws of mechanics the same
in every inertial frame”



Relativities

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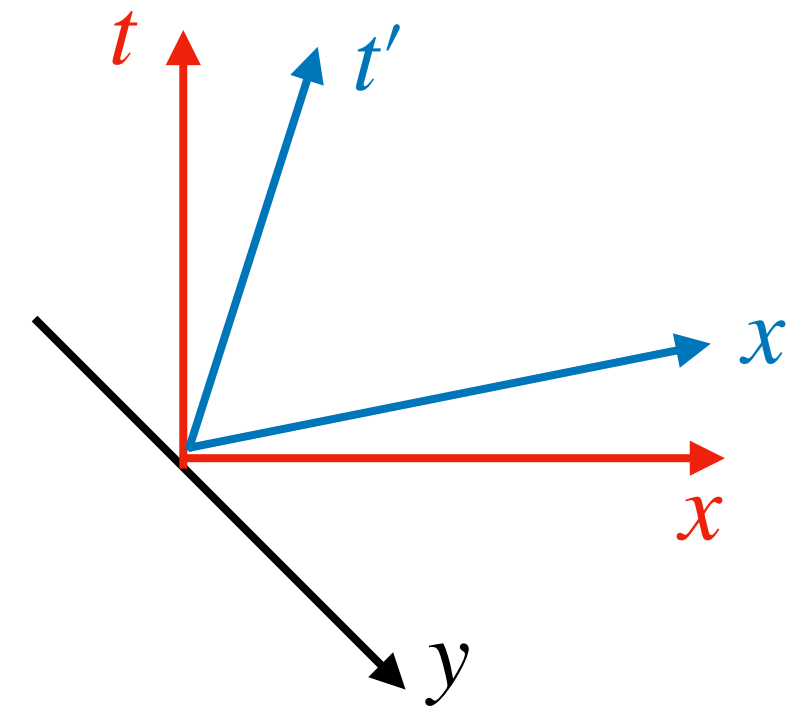
frame-dependent decomposition of space

$$\frac{1}{c} \neq 0$$



Special

“All the laws of physics (exc. Newt. gravity) the same in every inertial frame”

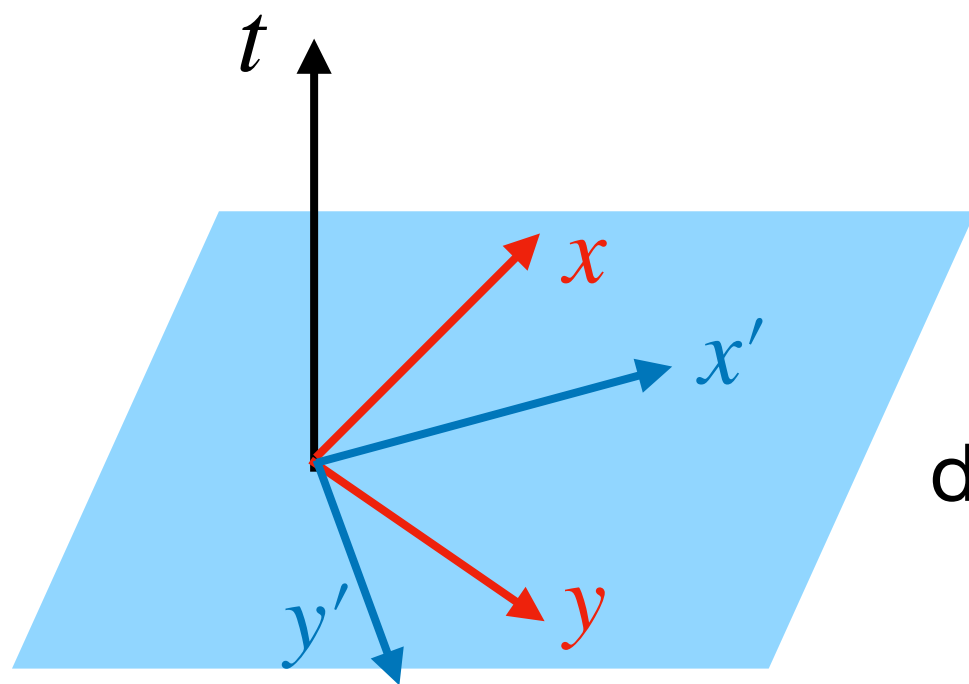


frame-dependent decomposition of spacetime into space and time

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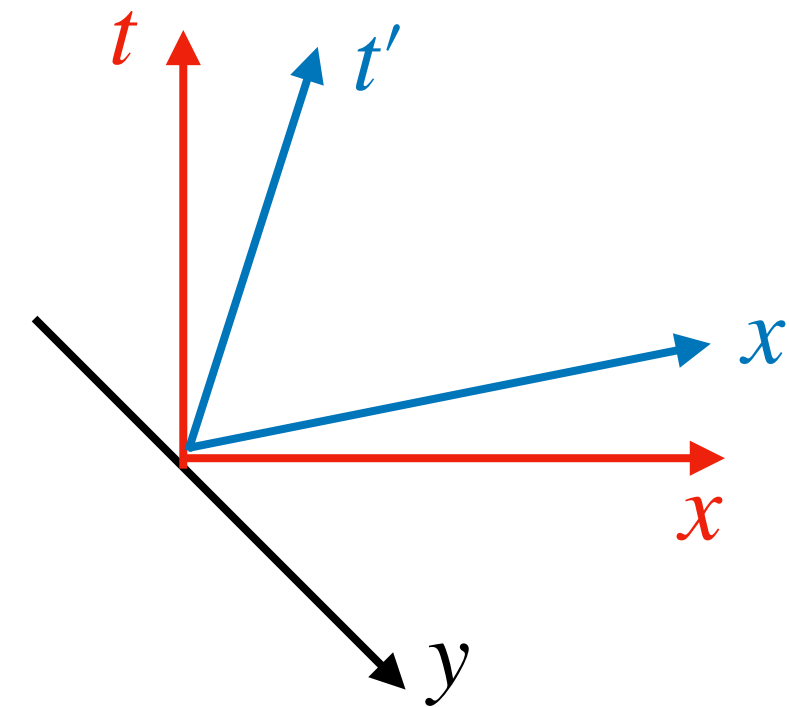
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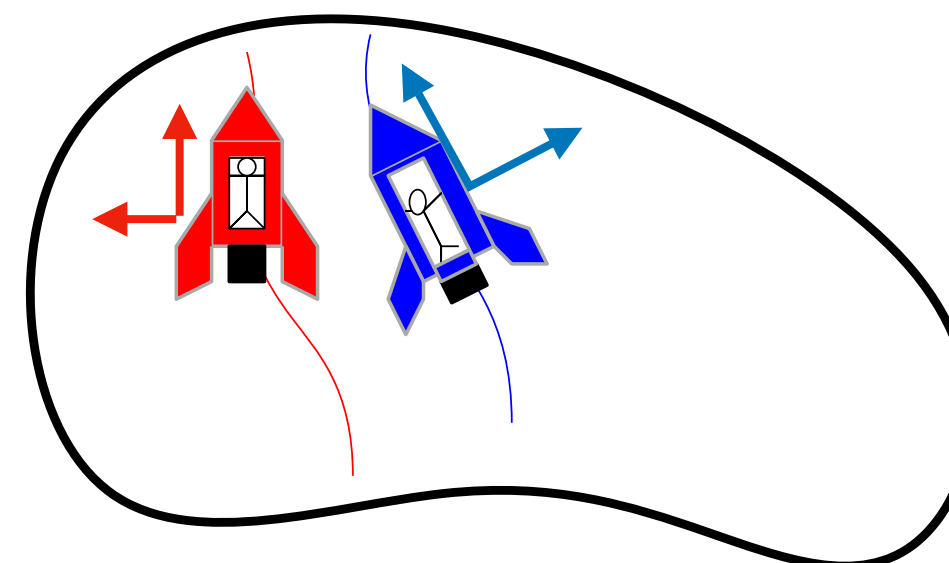
frame-dependent decomposition of spacetime into space and time

$$G \neq 0$$



General

“All the laws of physics the same in every frame”

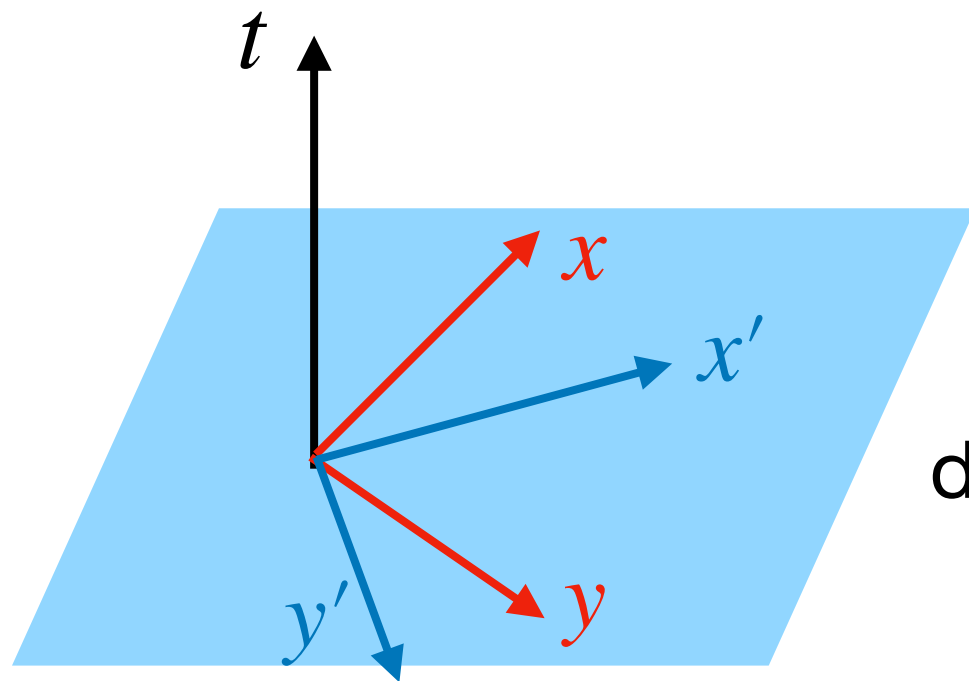


frame-dependent local decomposition of spacetime into space and time

Relativities

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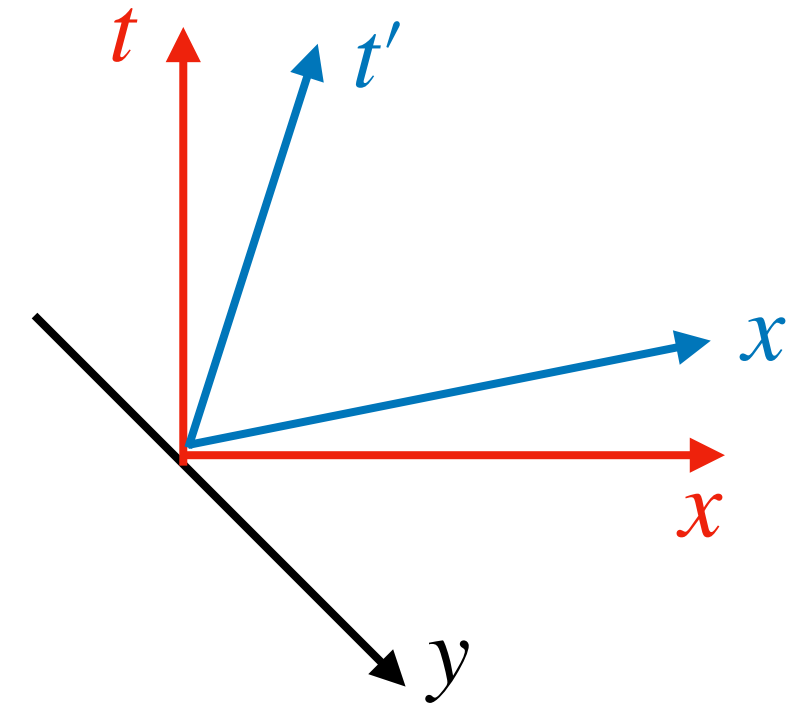
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frame-dependent decomposition of spacetime into space and time

$$G \neq 0$$



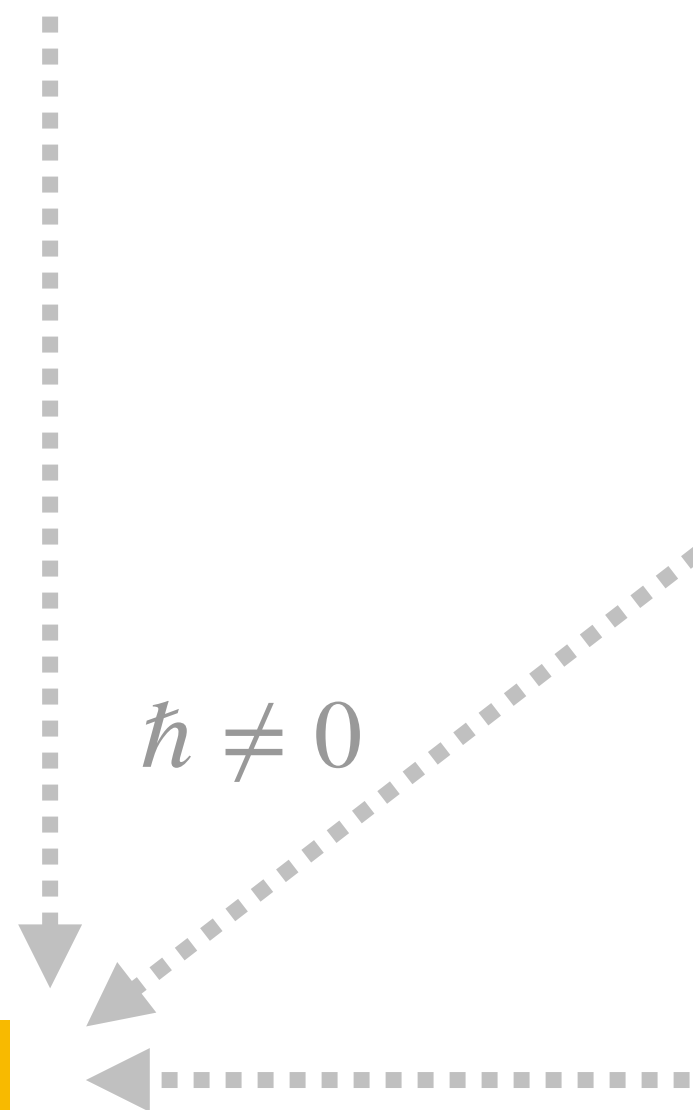
Quantum?

“All the ... laws of ... the same in every ... QRF”



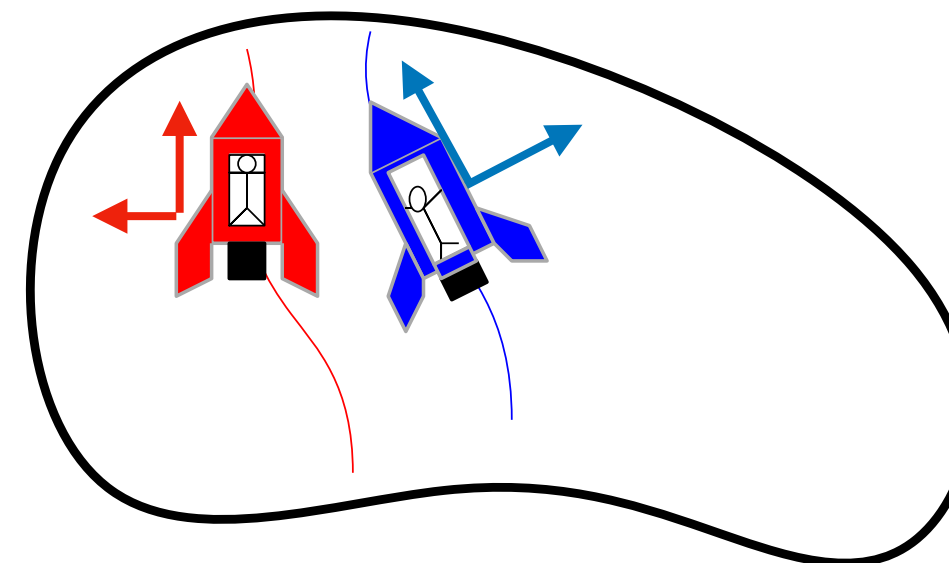
QRF-dependent decomposition of

$$\hbar \neq 0$$



General

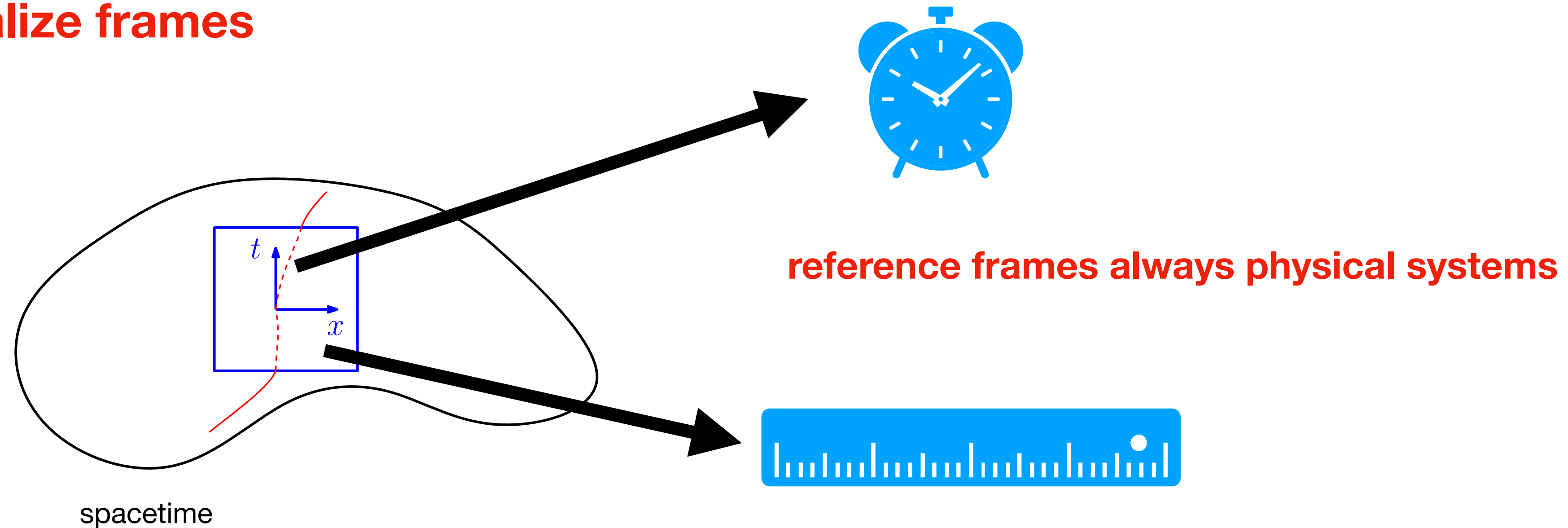
“All the laws of physics the same in every frame”



frame-dependent local decomposition of spacetime into space and time

Quantum reference frames

internalize frames



- no background structure

- universality of QT (ext. of Heisenberg cut)

⇒

RF subject to QM itself

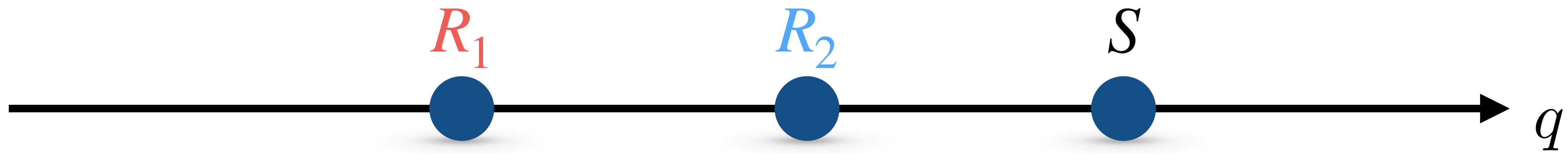
“RFs in relative superposition”

Why care?

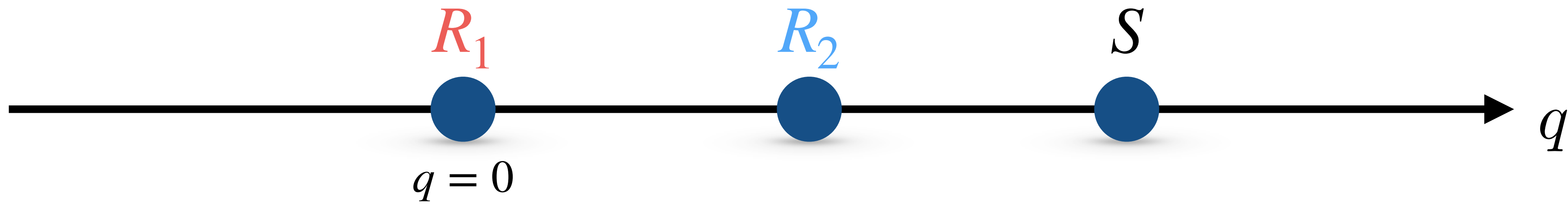
- **Foundational interest**
classical frame relations \Leftrightarrow classical spacetime structure

 \Rightarrow quantum frame relations \Leftrightarrow quantum spacetime structure?
- **systems with gauge symmetry (gauge-inv. descriptions implicitly invoke internal frames)**
- **gravity: no background frame**
- **quantum info: agents may not share a common external lab frame**

Intuition: spatial QRFs in 1D

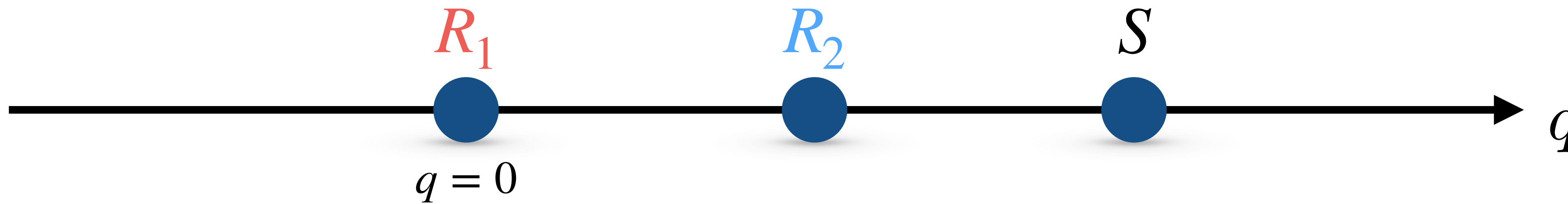


Intuition: spatial QRFs in 1D



aim: “jump” into perspective of particle R_1 , defining origin, and describe R_2S relative to it

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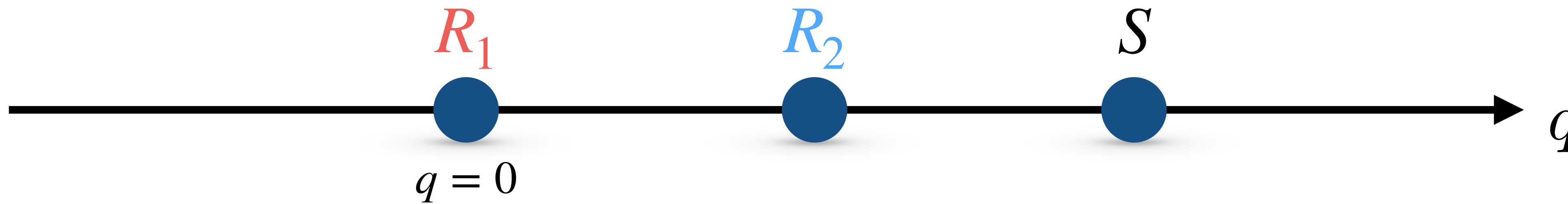


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$$|q_1\rangle_{R_2} \otimes |x\rangle_S$$

R_1 perspective

Intuition: spatial QRFs in 1D



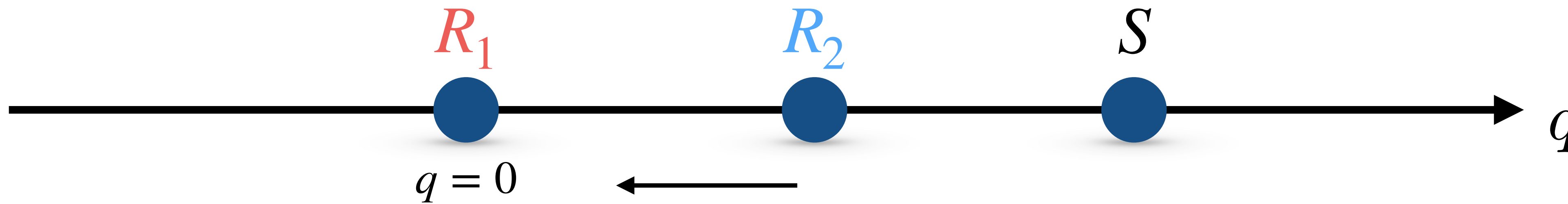
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how will R_2 “see” the same configuration?

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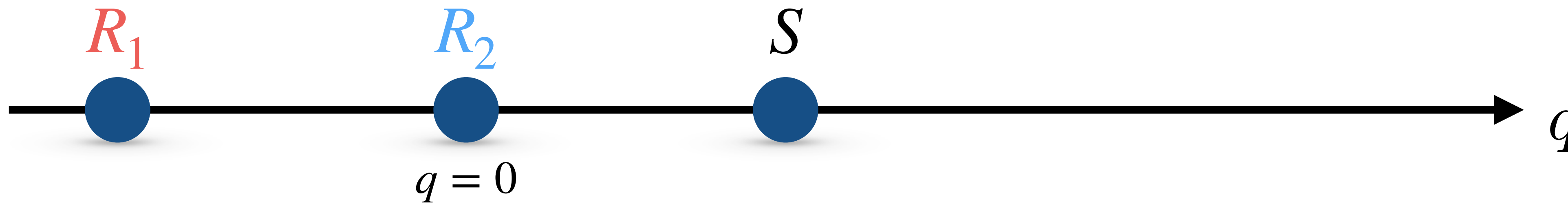
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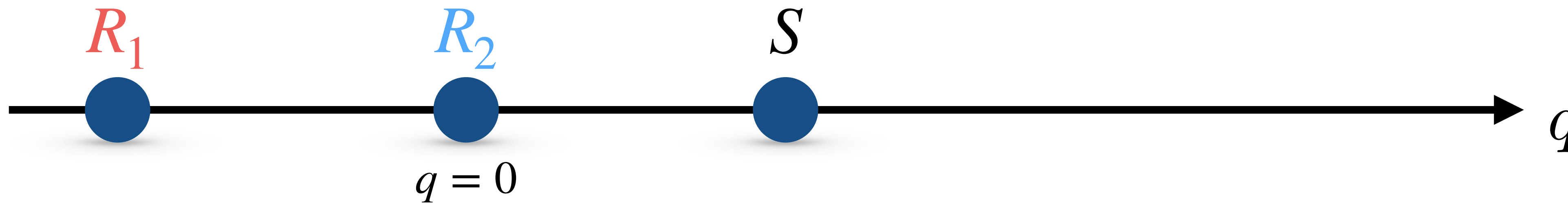
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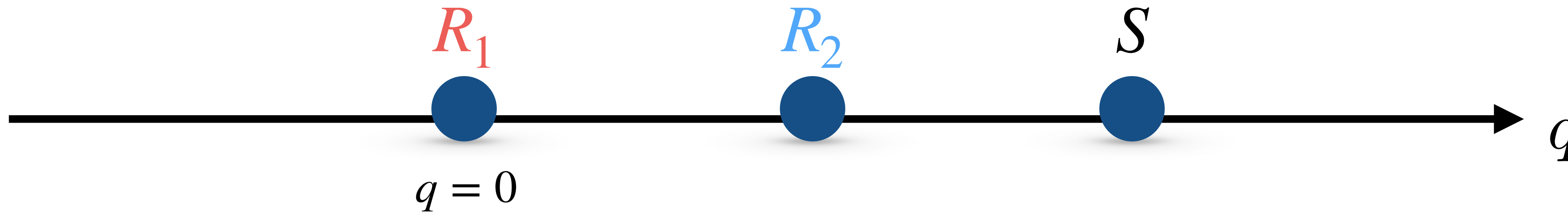


$$|-q_1\rangle_{R_1} \otimes |x - q_1\rangle_S$$

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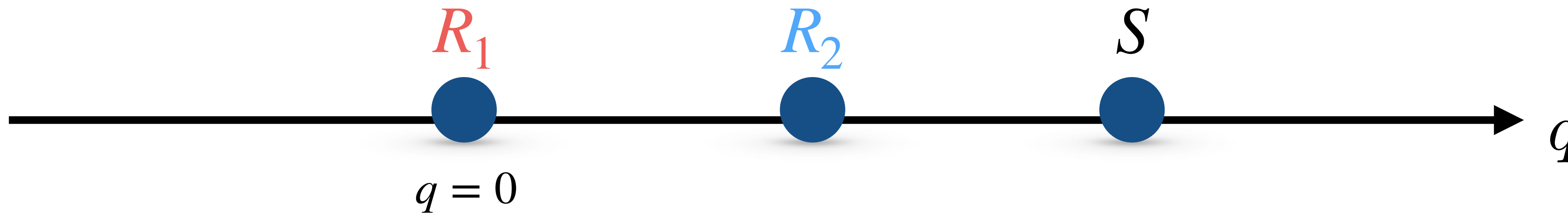


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$$\begin{array}{ccc} |q_1\rangle_{R_2} \otimes |x\rangle_S + |q_2\rangle_{R_2} \otimes |x\rangle_S & \longrightarrow & | -q_1\rangle_{R_1} \otimes |x - q_1\rangle_S \\ R_1 \text{ perspective} & & R_2 \text{ perspective} \end{array}$$

how will R_2 “see” the same configuration? \Rightarrow what about superpositions?

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R_1 perspective



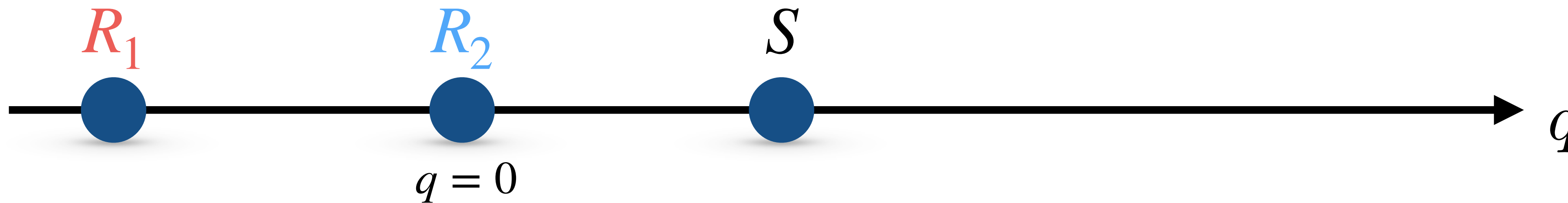
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R_2 perspective

assuming linearity

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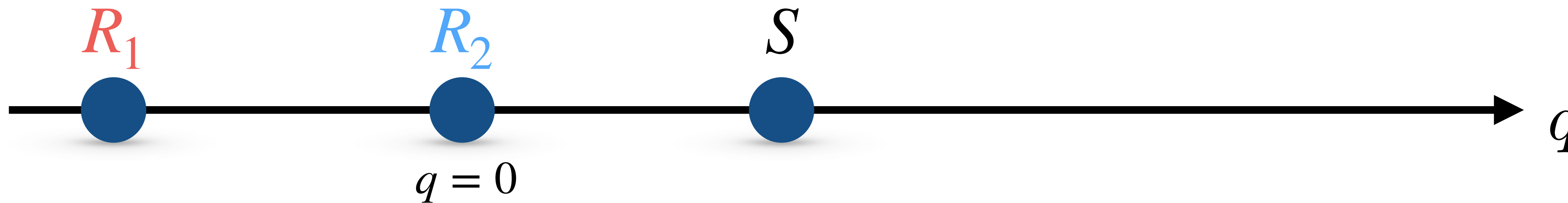
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QRF transformation a conditional unitary:

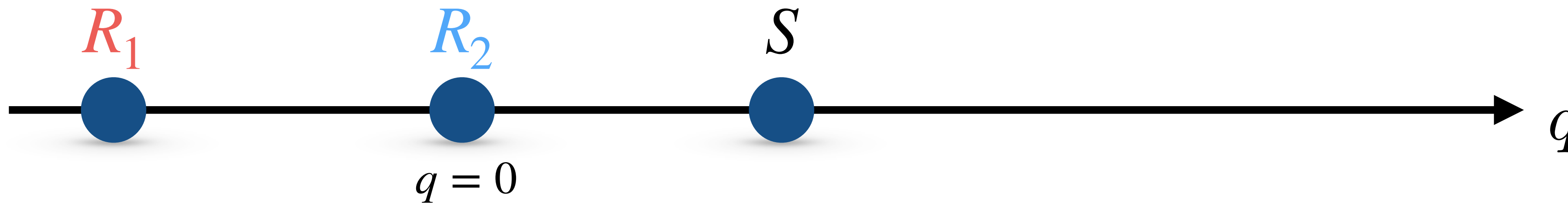
$$V_{R_1 \rightarrow R_2} = \mathbb{F}_{12} \int dq | -q \rangle \langle q |_{R_2} \otimes U_S(-q)$$

swaps particles R_2 and R_1

how will R_2 “see” the same configuration?

[Giacomini et al Nat. Comm. '19]

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[Giacomini et al Nat. Comm. '19]

Example illustrates: superposition and entanglement of subsystem S QRF relative

The story more generally

RFs and symmetries

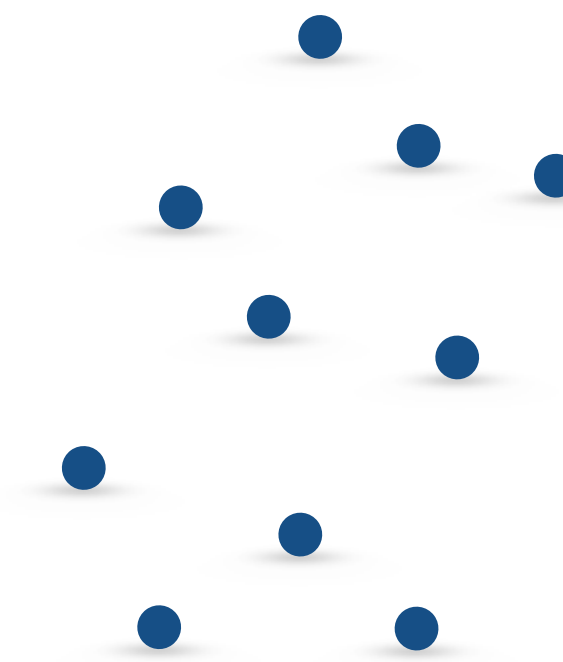
Premise:

System \mathcal{S} subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when \mathcal{S} considered in isolation

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pair (G, \mathcal{S}) could be, e.g.:

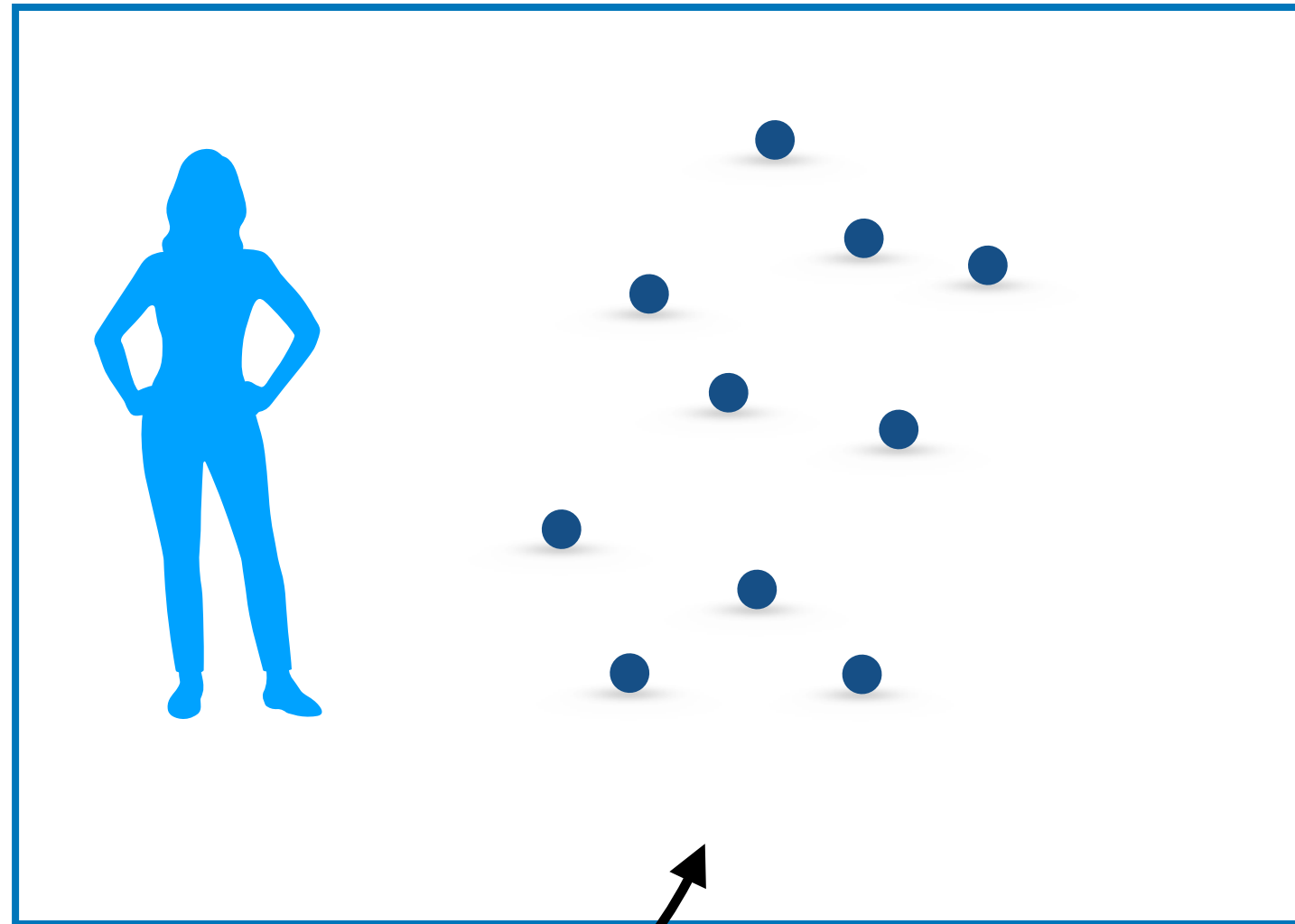
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- diffeos + all dynamical fields in spacetime
- ...



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external frame

quantum information/foundations:
lab frame

gravity:
fictitious (or edge modes)

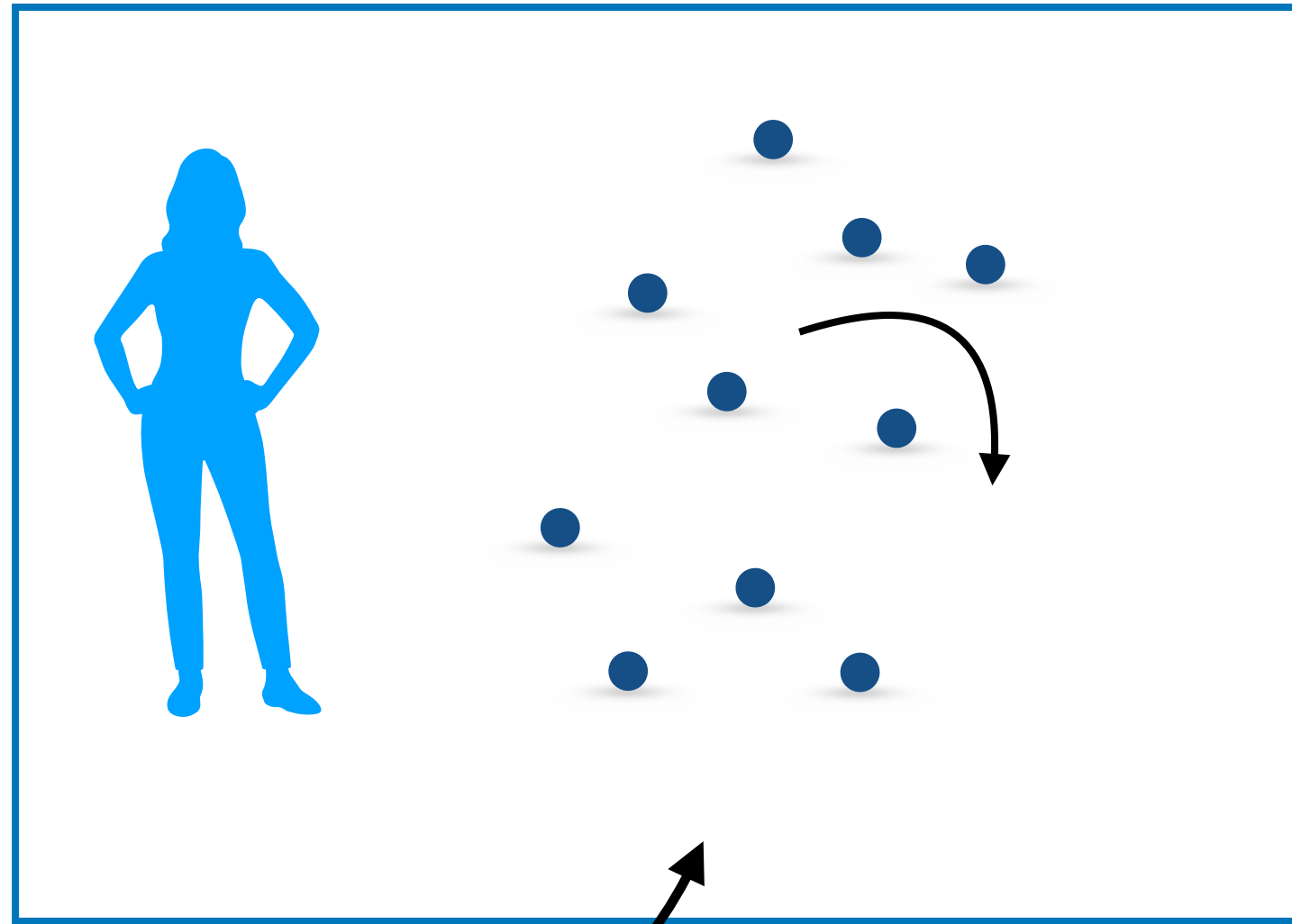
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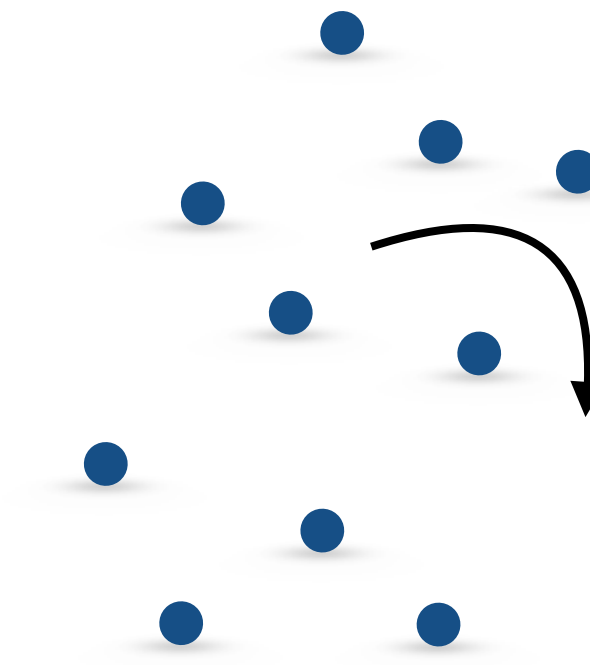
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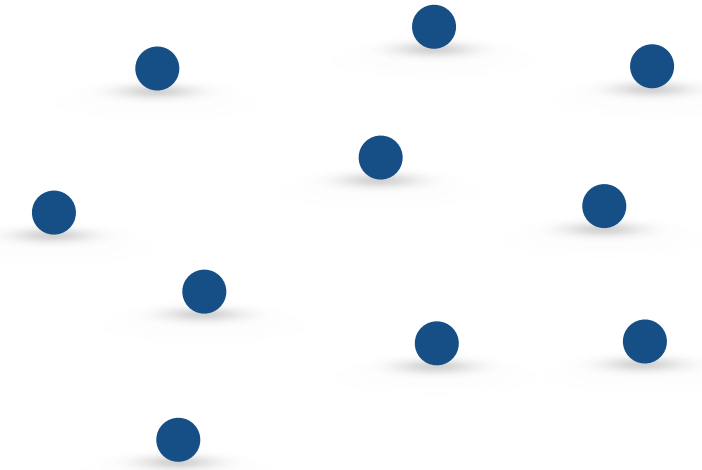
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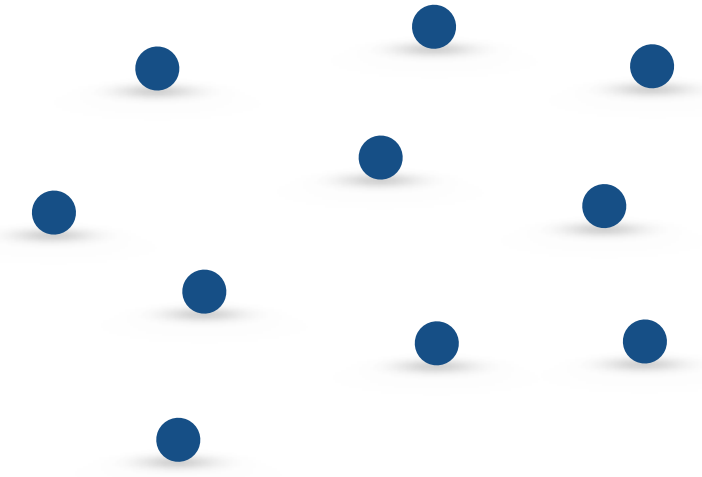
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\Rightarrow “gauge inv.” = external frame indep.

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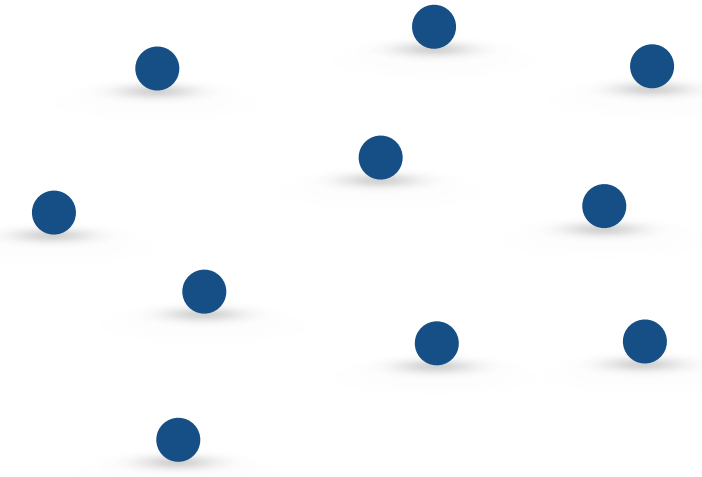
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quantum information/foundations:
change of ext. lab frame

gravity:
change of background coordinates (diffeo)
 \Rightarrow change of fictitious background frame

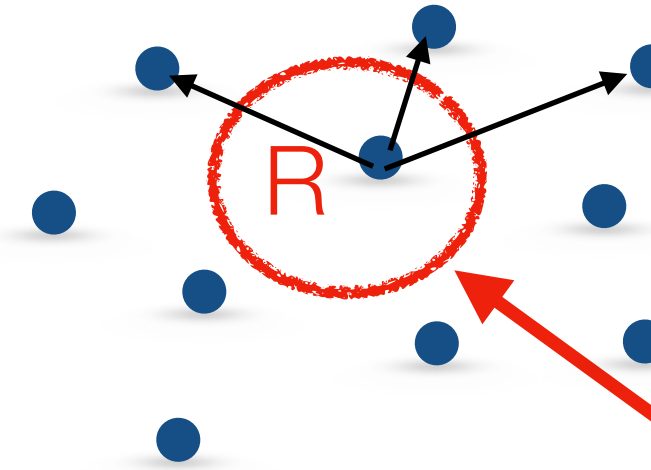
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Describe S relative to internal reference subsystem R

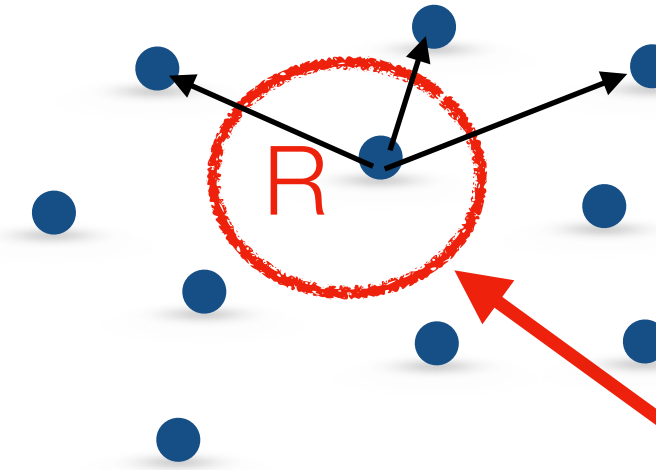
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**Describe S relative to internal
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should transform “nicely” (covariantly) under G

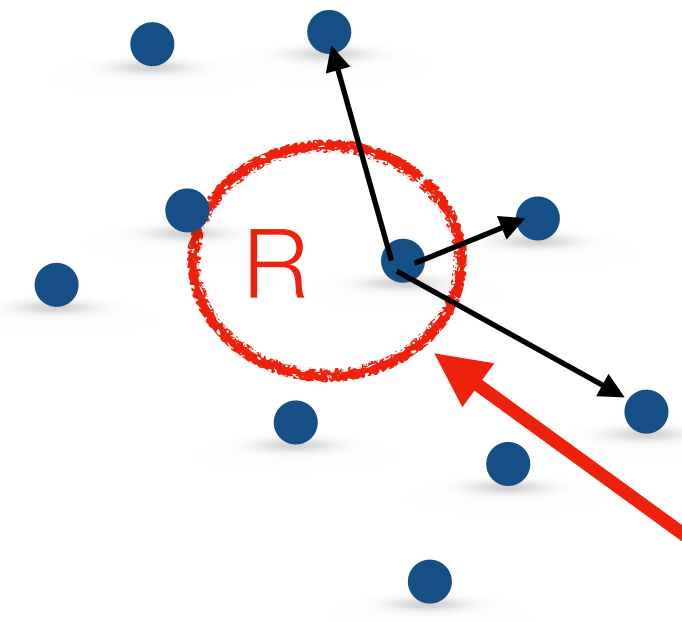
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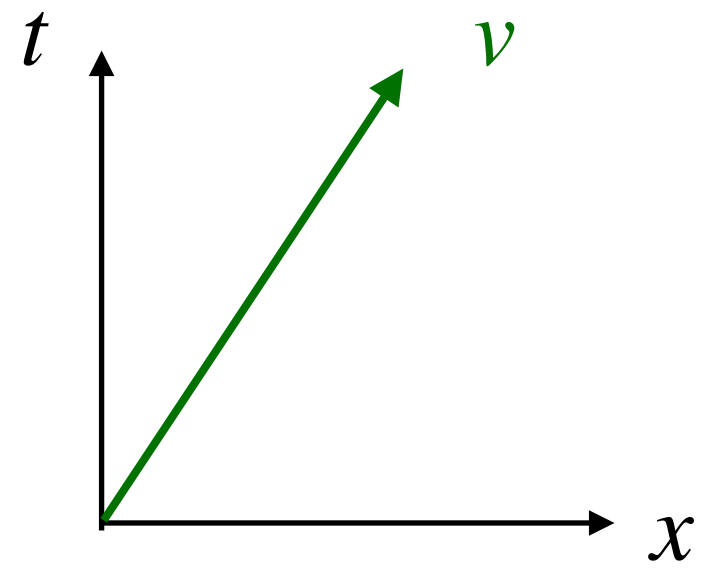
internally indistinguishable

Describe S relative to internal reference subsystem R

should transform “nicely” (covariantly) under G

ρ and $g \cdot \rho$ members of same relational equivalence class of states, different descriptions of same relational state

Example: Special relativity with internal frames



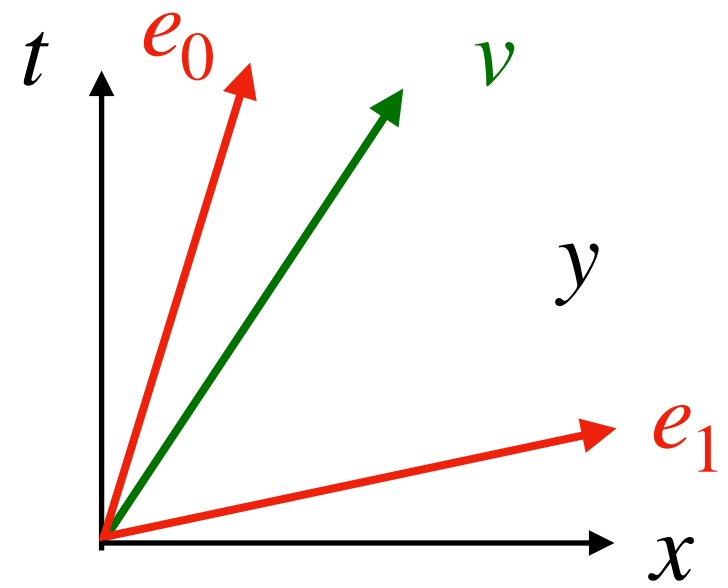
$$v^\mu \mapsto \Lambda^\mu{}_\nu v^\nu$$

$$\Lambda \in \text{SO}_+(3,1),$$

internally indistinguishable

fictitious/external coord. frame

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introduce internal frame (tetrad)

e_a^μ

$\mu = t, x, y, z$ spacetime index,

$a = 0, 1, 2, 3$

frame index

fictitious/external coord. frame

frame orientations

\Rightarrow group acts on itself since

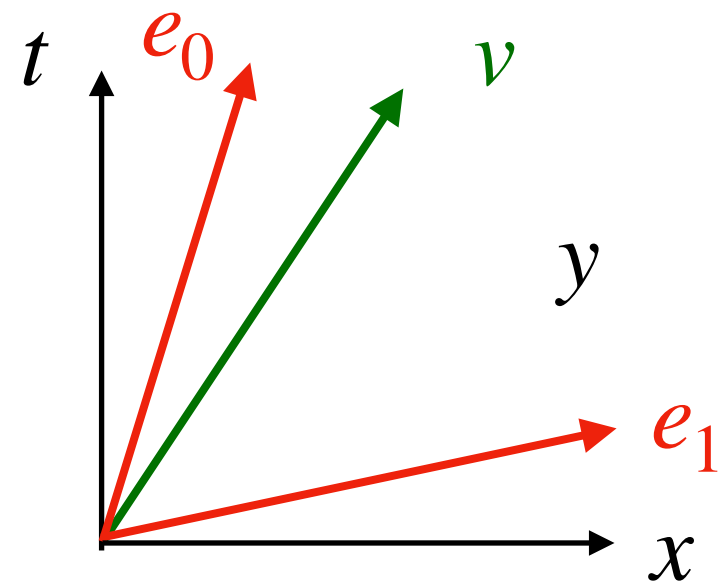
$$\eta_{ab} = e_a^\mu e_b^\nu \eta_{\mu\nu}$$

\Rightarrow

$$e_a^\mu \in \text{SO}_+(3,1)$$

group valued frame orientations

Example: Special relativity with internal frames



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• “gauge transformations”:

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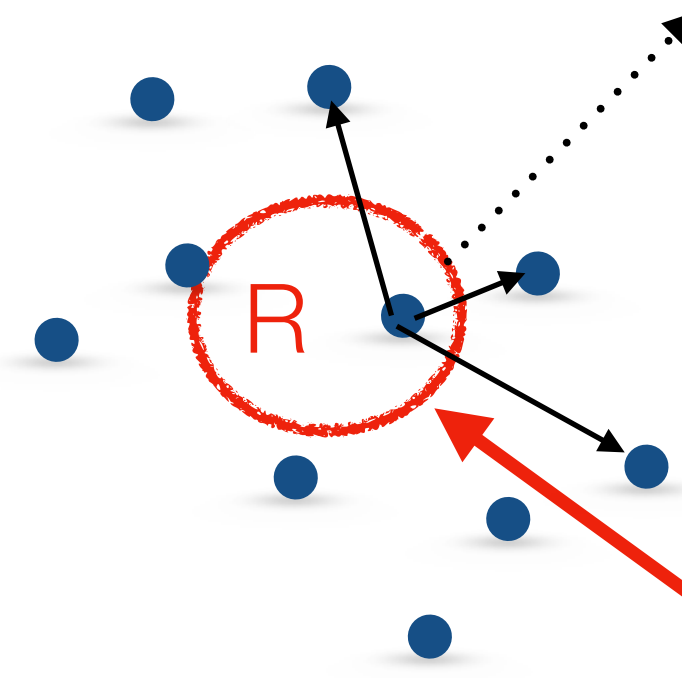
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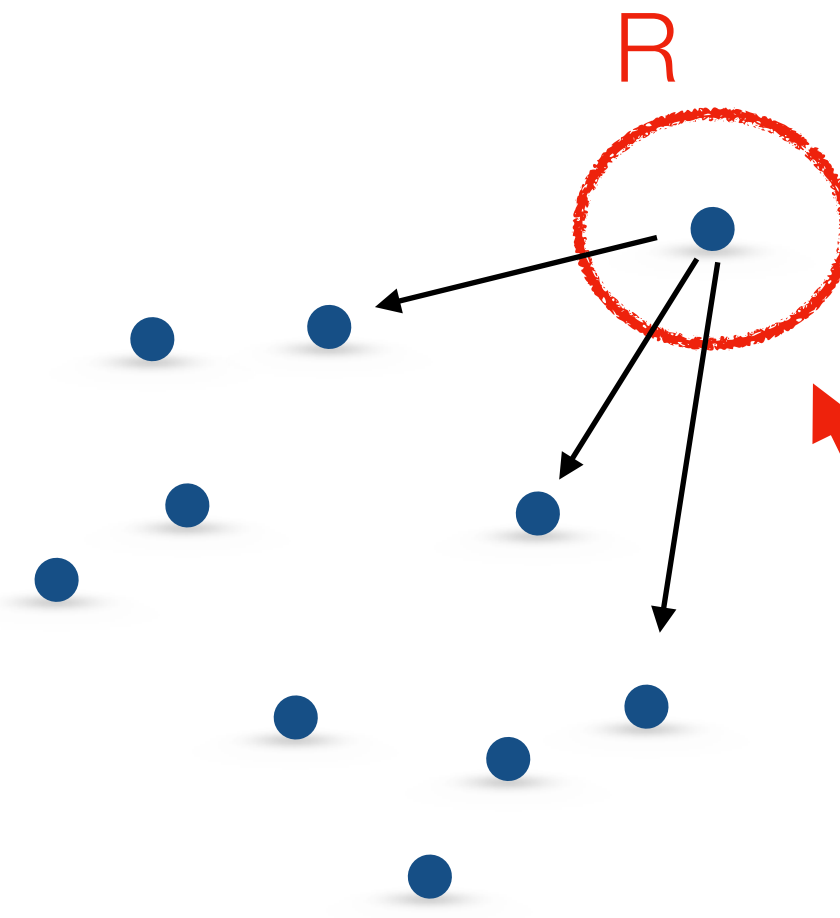
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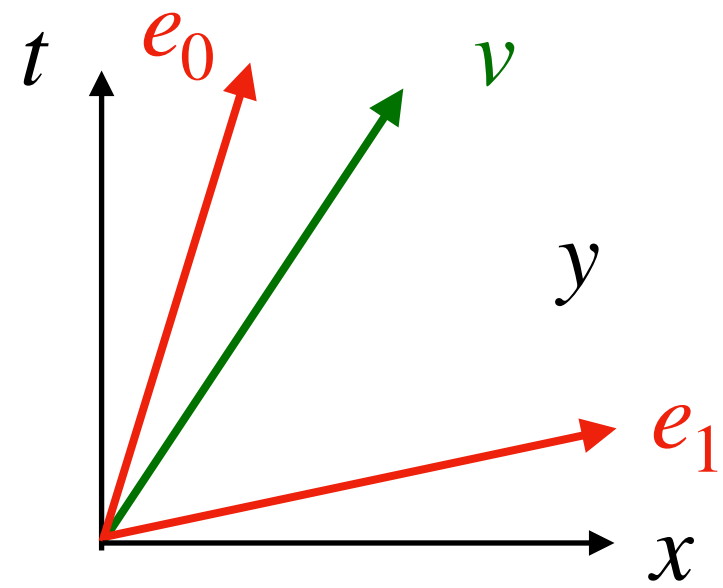
internally distinguishable
(relations changed)

⇒ **internal frame reorientation**
("symmetry")

**Describe S relative to internal
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**Interested in internally
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fictitious/external coord. frame

frame orientations

2 indices, 2 **commuting** group actions:

- “gauge transformations”:
- “symmetries” (frame reorientations):

$$\Lambda^\mu{}_\nu e_a^\nu \quad \Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$$

$$\Lambda_a{}^b e_b^\mu \quad \Lambda_a{}^b \in \text{SO}_+(3,1)$$

only acts on frame

\Rightarrow group acts on itself since

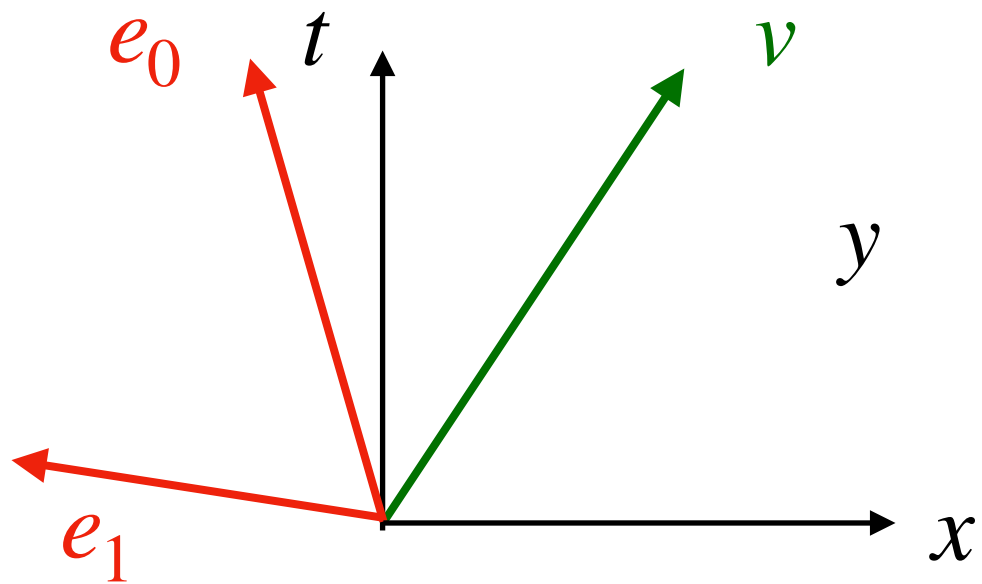
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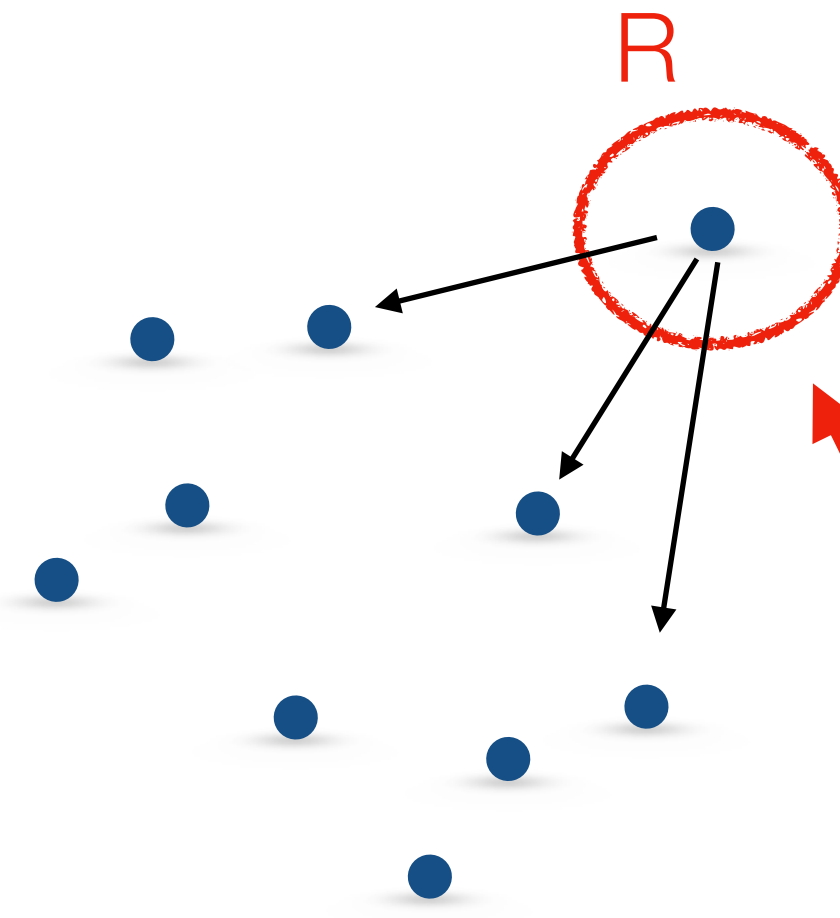
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group valued frame orientations

2 ways of “jumping into a RF perspective”

Premise:

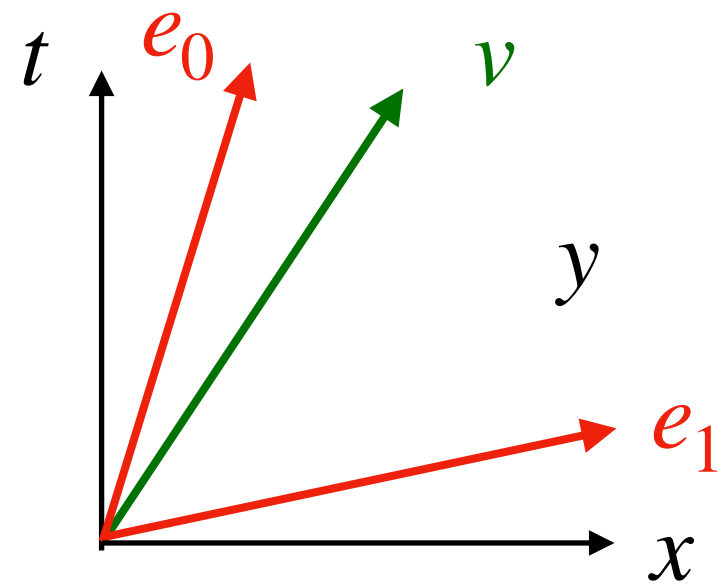
System S subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when S considered in isolation



Describe S relative to internal reference subsystem R

1. relational observables relative to R (gauge inv.)

Example: Special relativity with internal frames



$$v^\mu \mapsto \Lambda^\mu{}_\nu v^\nu$$

$$\Lambda \in \text{SO}_+(3,1),$$

internally indistinguishable

introduce internal frame (tetrad)

e_a^μ

$\mu = t, x, y, z$ spacetime index,

$a = 0, 1, 2, 3$

frame index

fictitious/external coord. frame

frame orientations

2 indices, 2 **commuting** group actions:

- “gauge transformations”:

$$\Lambda^\mu{}_\nu e_a^\nu \quad \Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$$

- “symmetries” (frame reorientations):

$$\Lambda_a{}^b e_b^\mu \quad \Lambda_a{}^b \in \text{SO}_+(3,1)$$

only acts on frame

⇒ group acts on itself since

$$\eta_{ab} = e_a^\mu e_b^\nu \eta_{\mu\nu}$$

⇒

$$e_a^\mu \in \text{SO}_+(3,1)$$

group valued frame orientations

⇒ “gauge-invariant” description of v :

$$v_a = (v, e_a) = \eta_{\mu\nu} v^\mu e_a^\nu$$

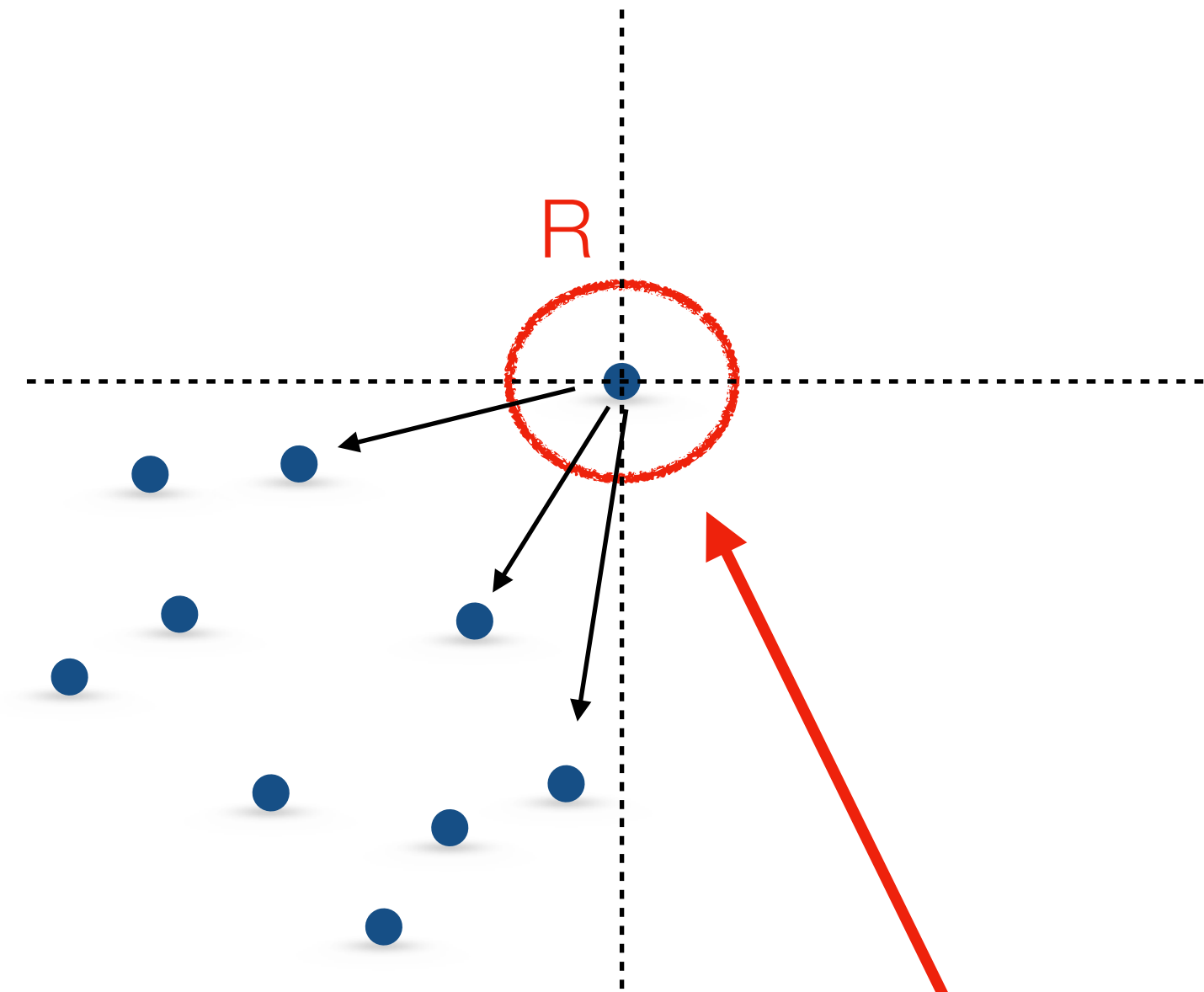
“relational/frame dressed observables”

(describes v relative to frame)

2 ways of “jumping into a RF perspective”

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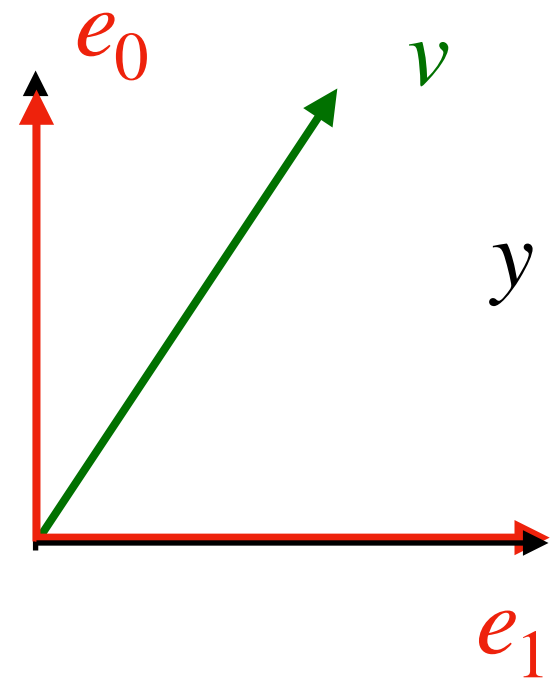
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Describe S relative to internal reference subsystem R

1. relational observables relative to R (gauge inv.)
2. put R into “origin” (gauge fix)

Example: Special relativity with internal frames



$$v^\mu \mapsto \Lambda^\mu{}_\nu v^\nu$$

$$\Lambda \in \text{SO}_+(3,1),$$

internally indistinguishable

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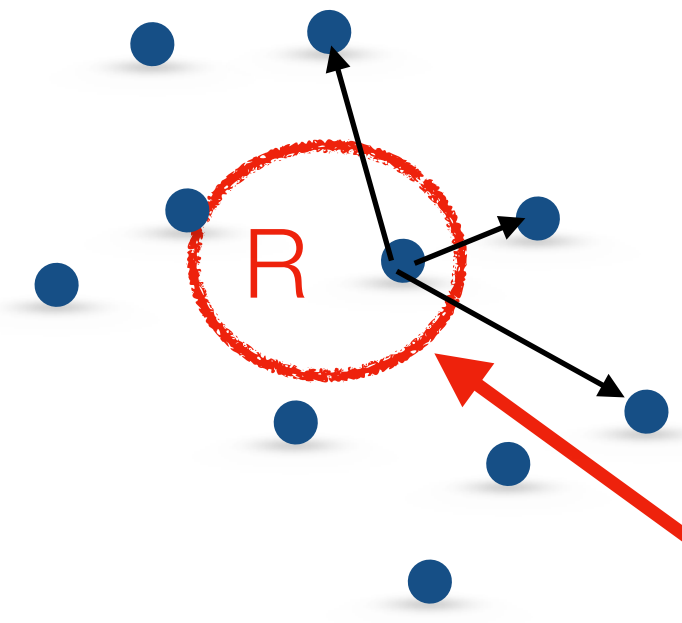
group valued frame orientations

⇒ gauge fix background frame to align with tetrad

The multiple choice problem

Premise:

System S subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when S considered in isolation



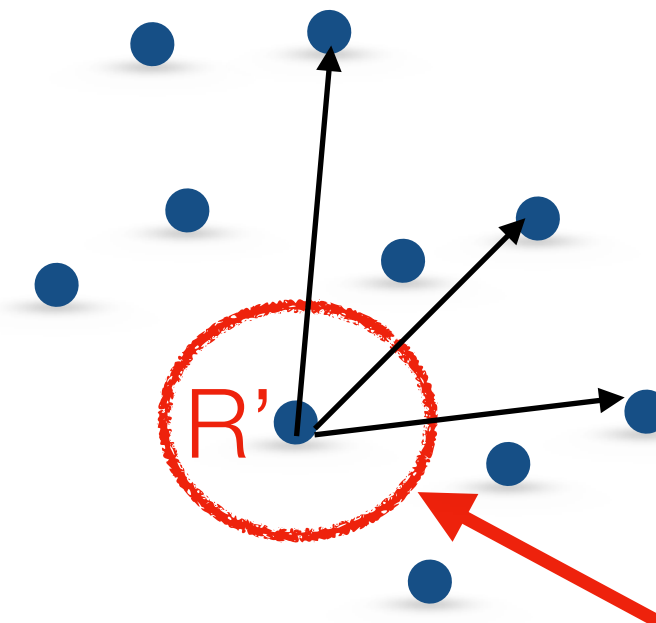
which frame to choose?

Describe S relative to internal reference subsystem R

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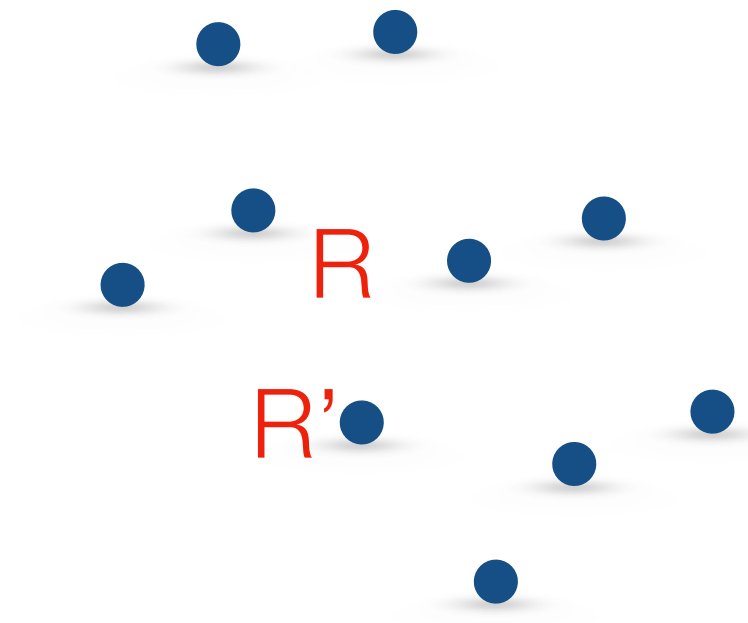
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Describe S relative to internal reference subsystem R'

2 ways of changing RF

Premise:

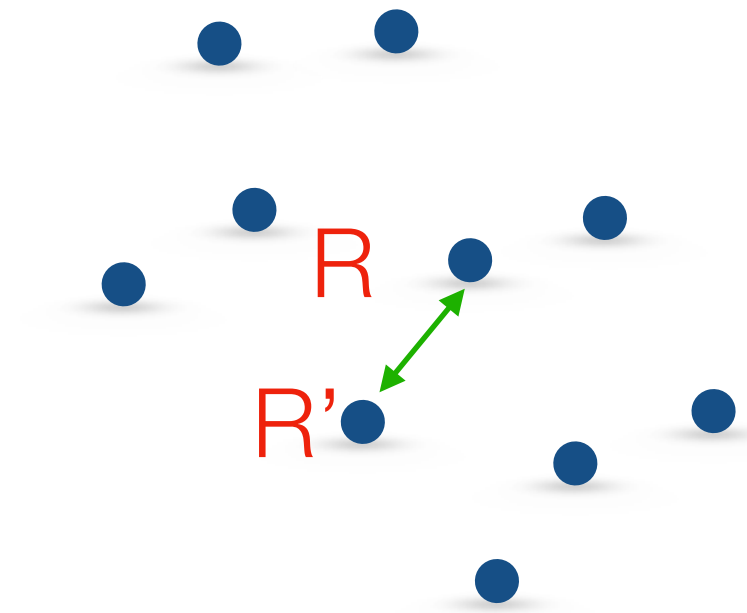
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2 ways of changing RF

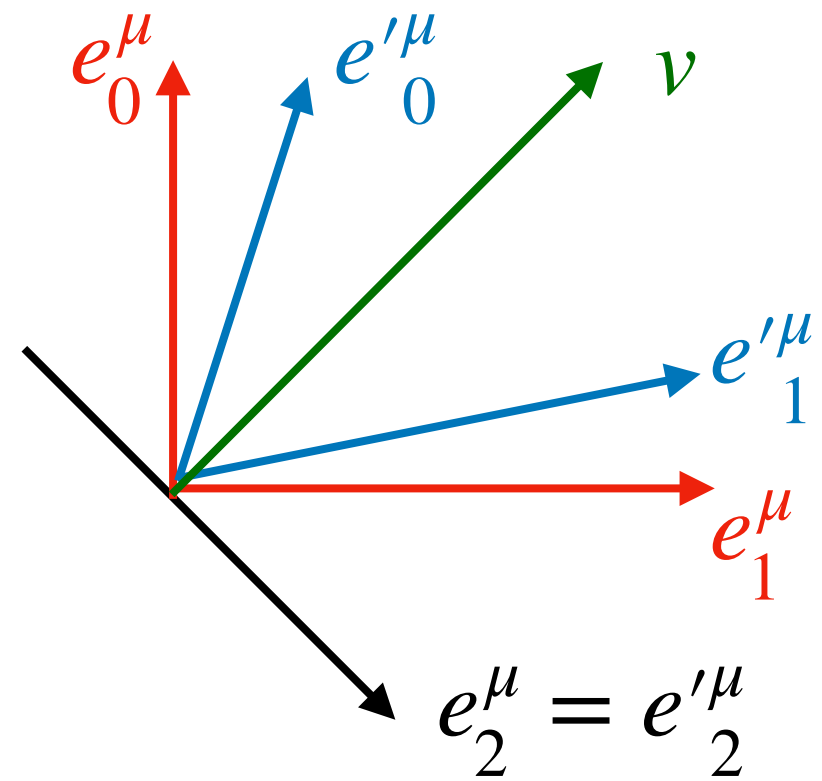
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1. relation-conditional reorientation

Warmup: Special relativity with internal frames



introduce second internal frame

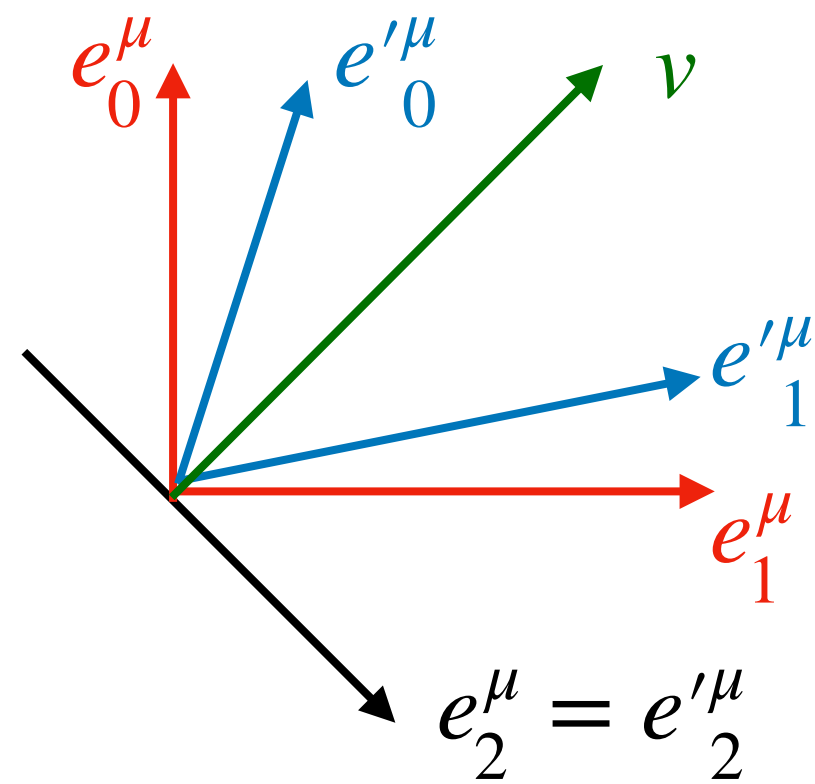
$e'_{a'}$

$$v_a = v^\mu \eta_{\mu\nu} e_a^\nu = v^\mu e'_{\mu a'} e_{\nu'}^{a'} e_a^\nu = v_{a'} \Lambda^{a'}_a$$

relational observable rel. to e

relational observable rel. to e'

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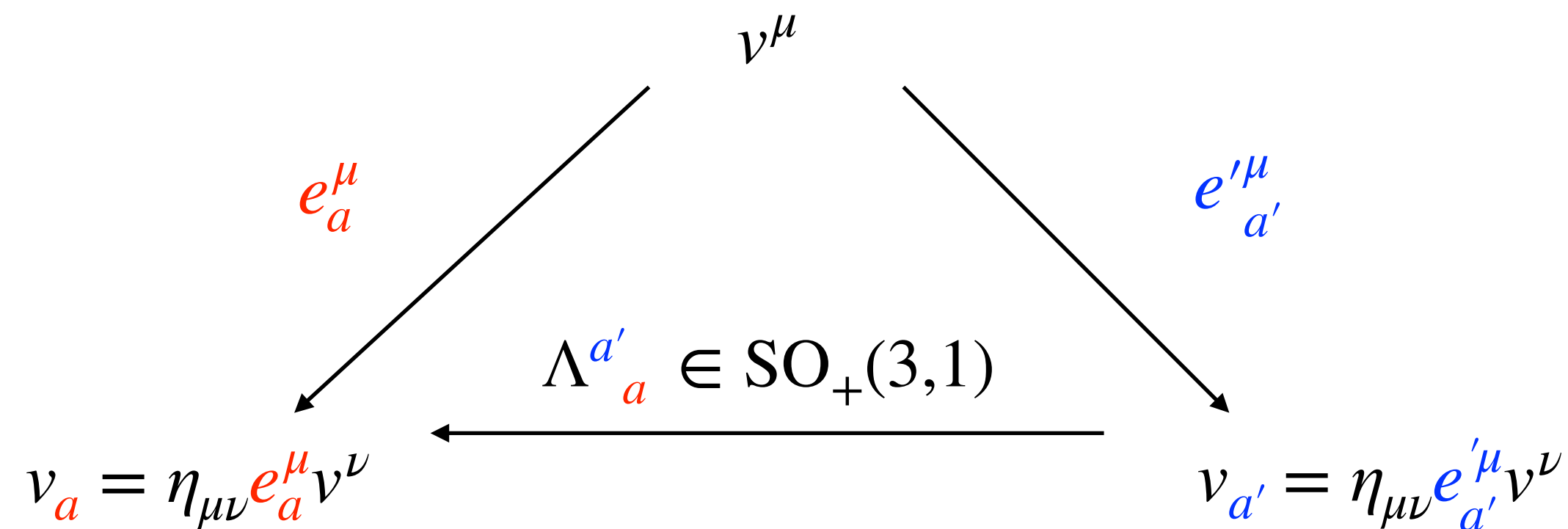
relational observable rel. to e

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symmetry induced RF transformation:

$$\Lambda^{a'}_a = e'^{a'}_\mu e_a^\mu \in \text{SO}_+(3,1)$$

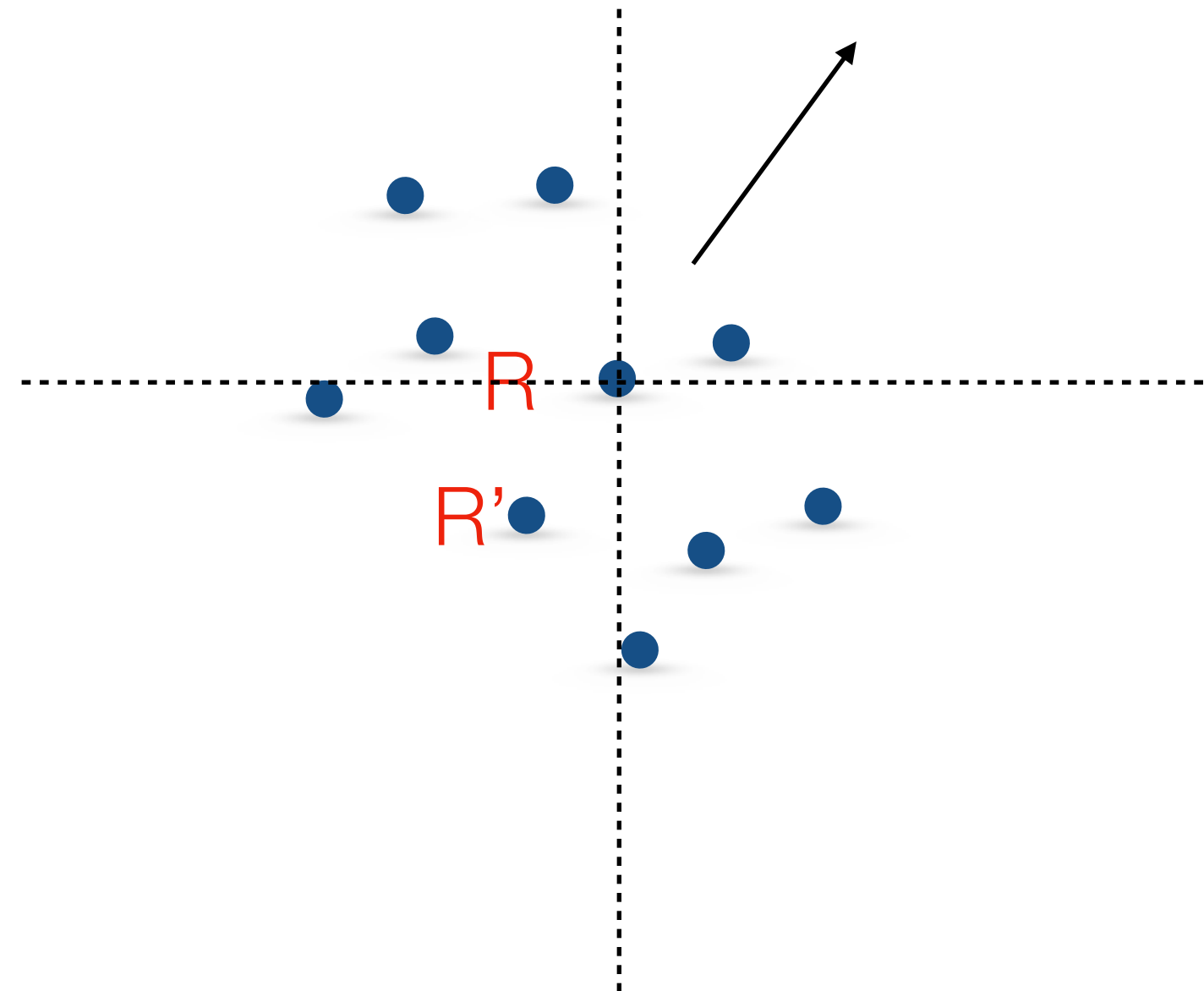
is relational observable describing 1st rel. to 2nd frame



2 ways of changing RF

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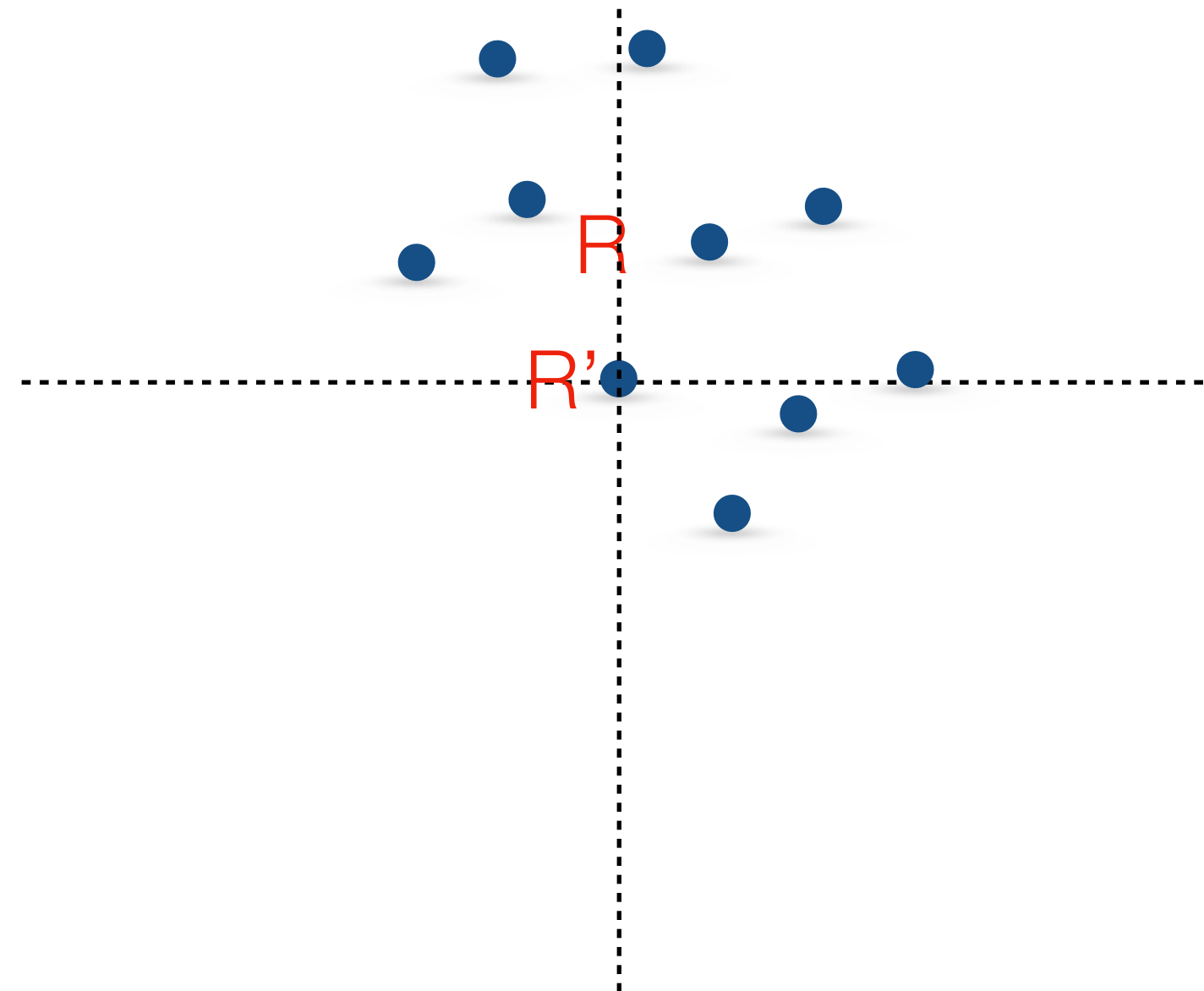


1. relation-conditional reorientation
2. relation-conditional gauge-transf.

2 ways of changing RF

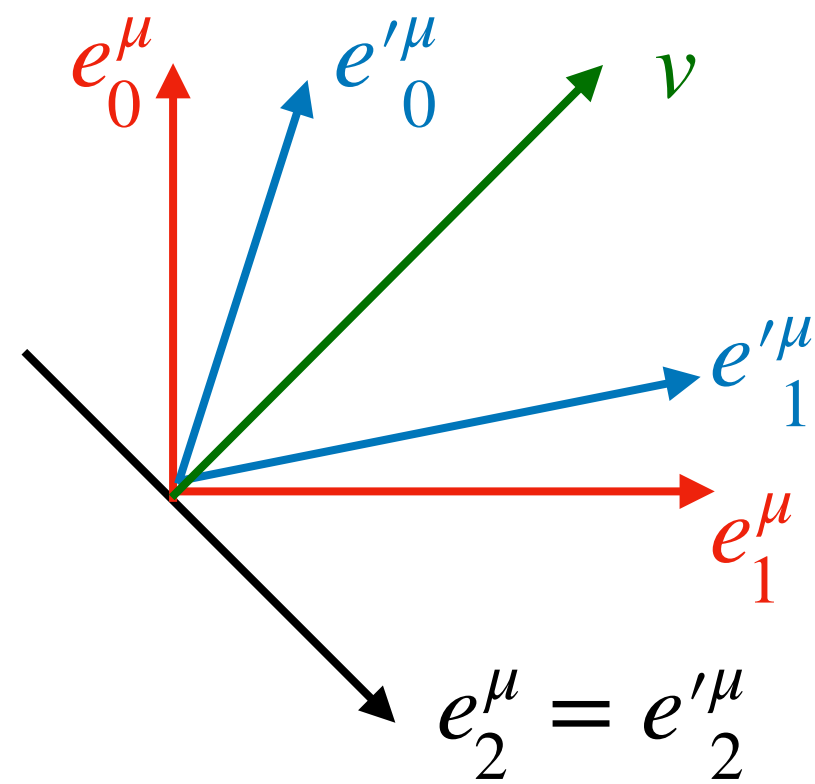
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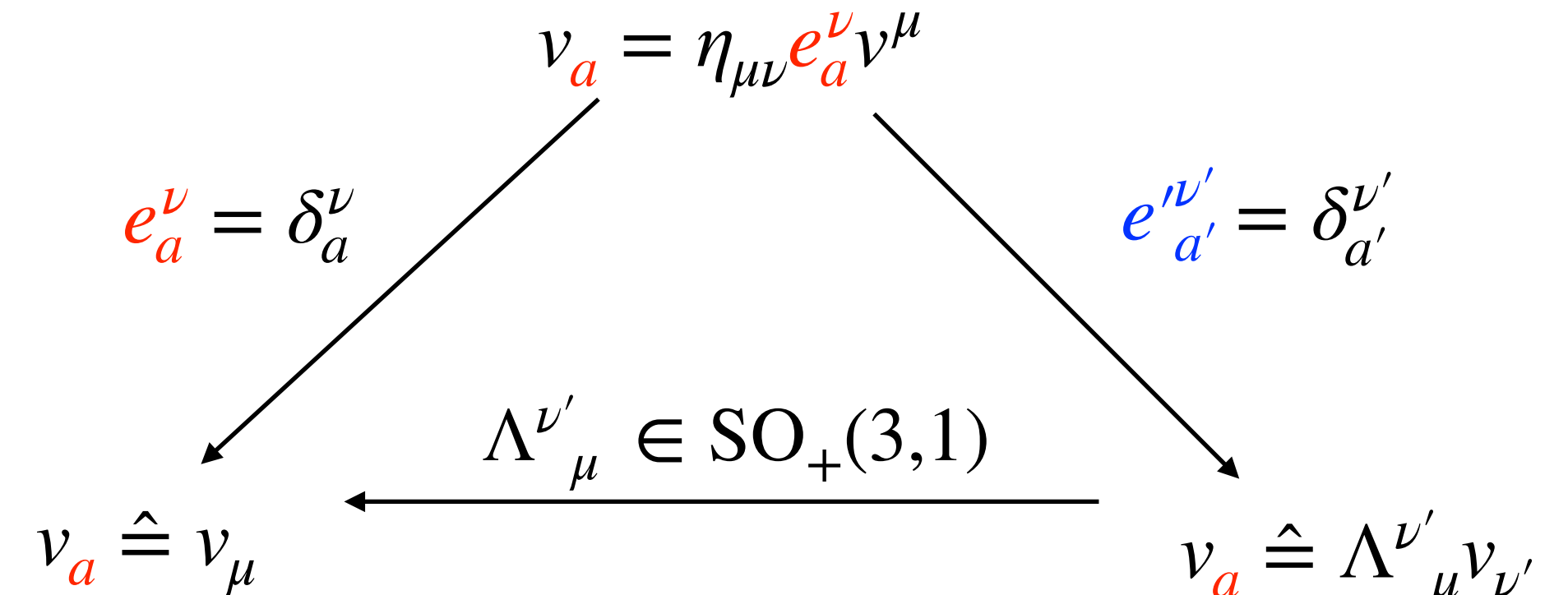
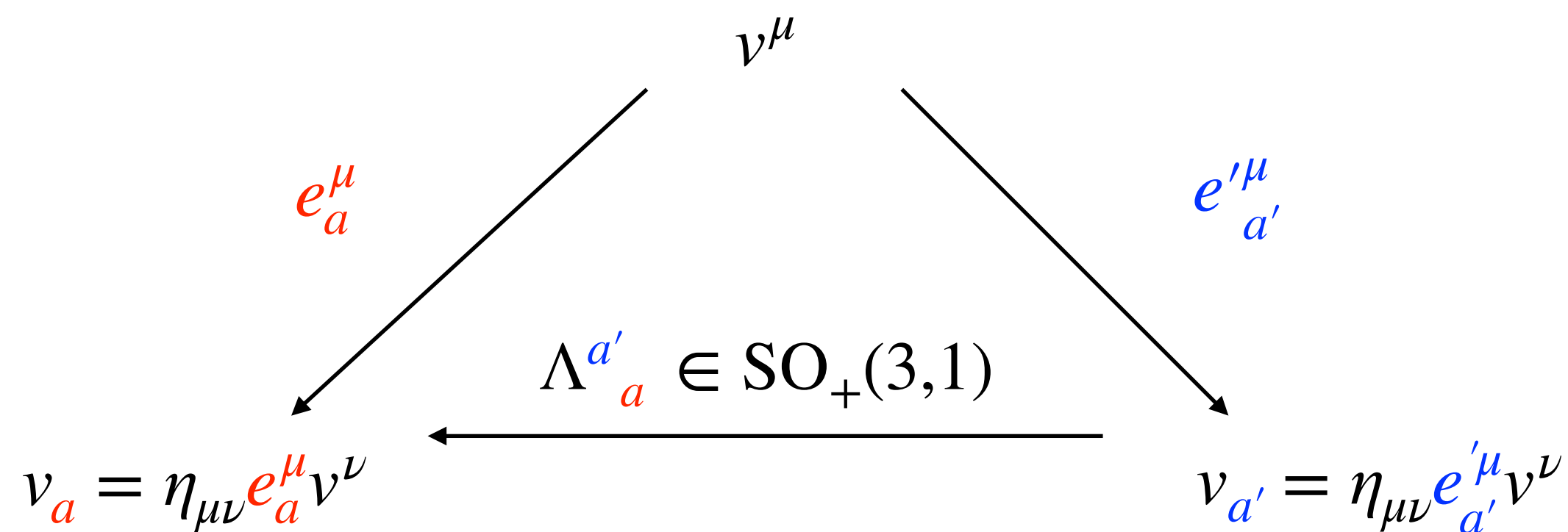
$$\Lambda^{a'}_a = e_{\mu'}^{a'} e_a^\mu \in \text{SO}_+(3,1)$$

is relational observable describing 1st rel. to 2nd frame

gauge induced RF transformation:

$$\Lambda^{\nu'}_\mu \in \text{SO}_+(3,1) \text{ looks the same as } \Lambda^{a'}_a$$

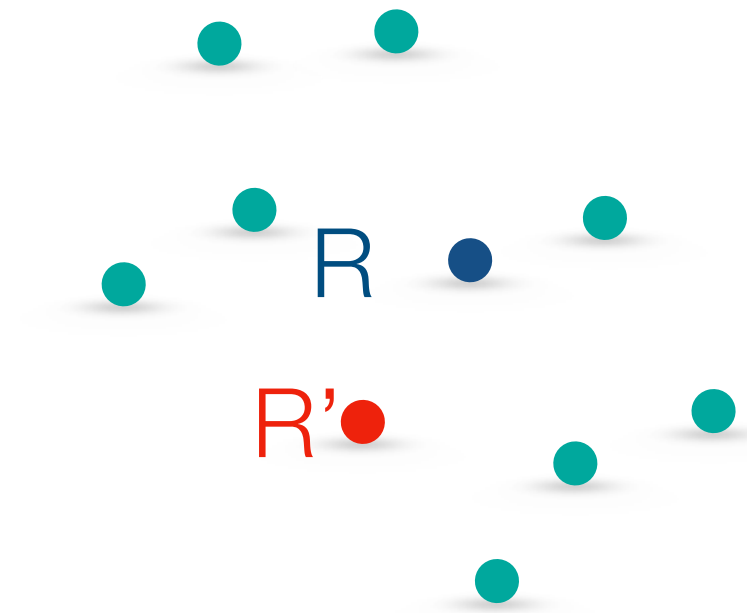
coordinate change via gauge fixings



kinematical vs. relational subsystems

Premise:

System \mathcal{S} subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when \mathcal{S} considered in isolation

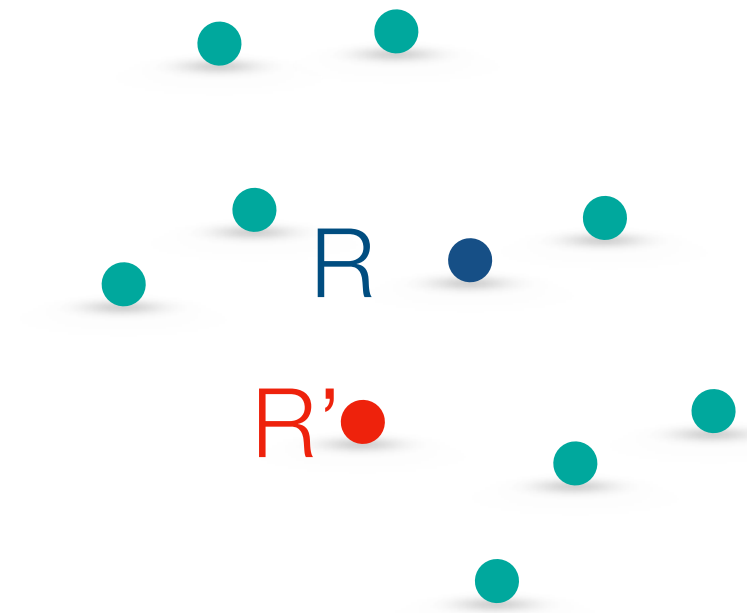


green balls: subsystem S'

kinematical vs. relational subsystems

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green balls: subsystem S'

1. kinematical and relational (gauge inv.) notion of subsystem **distinct**

kinematical vs. relational subsystems

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green balls: subsystem \mathcal{S}'



leaves description of \mathcal{S}' rel. to external frame invariant, but changes description relative to frame R

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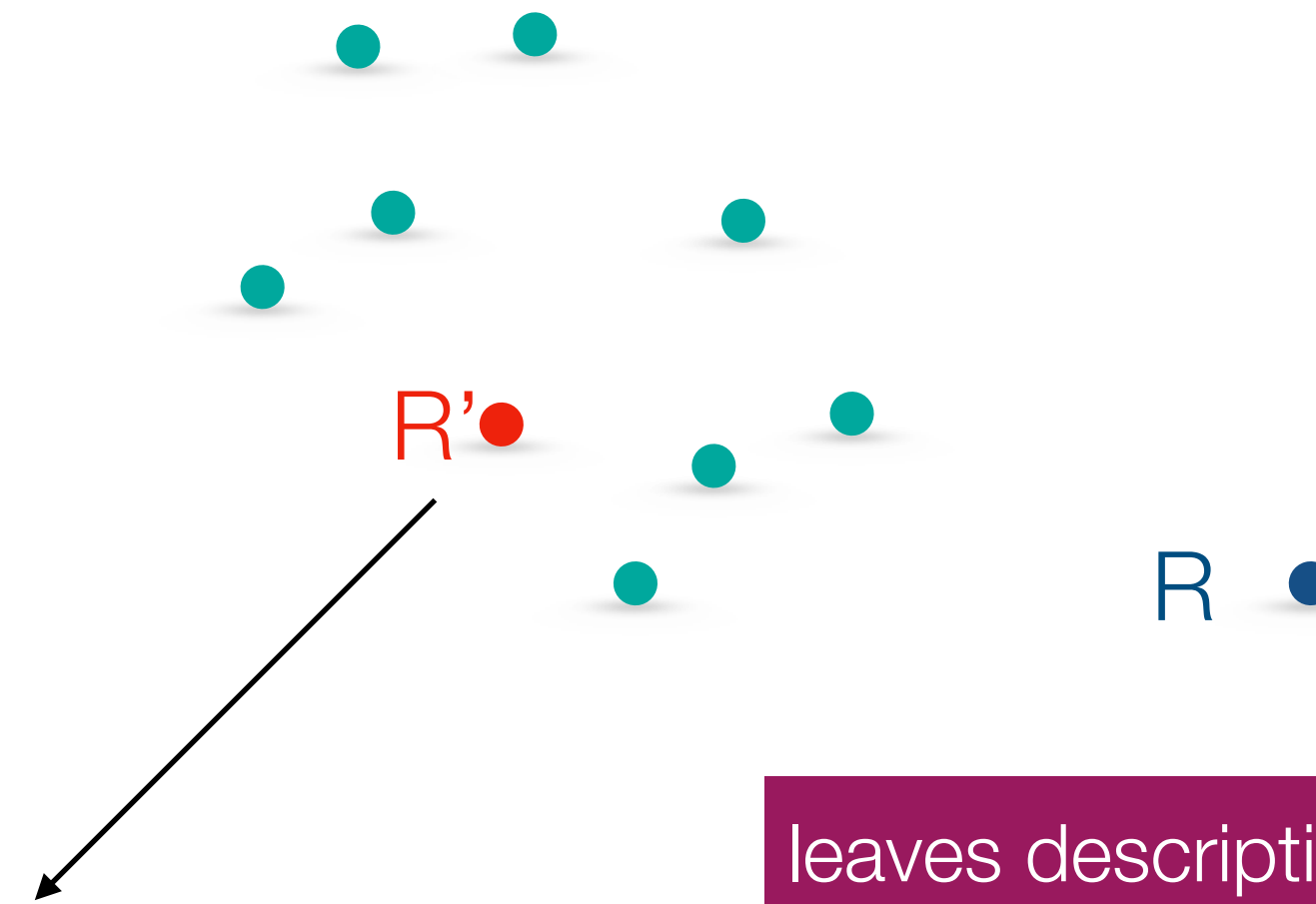
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kinematical vs. relational subsystems

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leaves description of S' relative to R invariant,
but changes it relative to R'

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Indistinguishable for all $g \in G$ when
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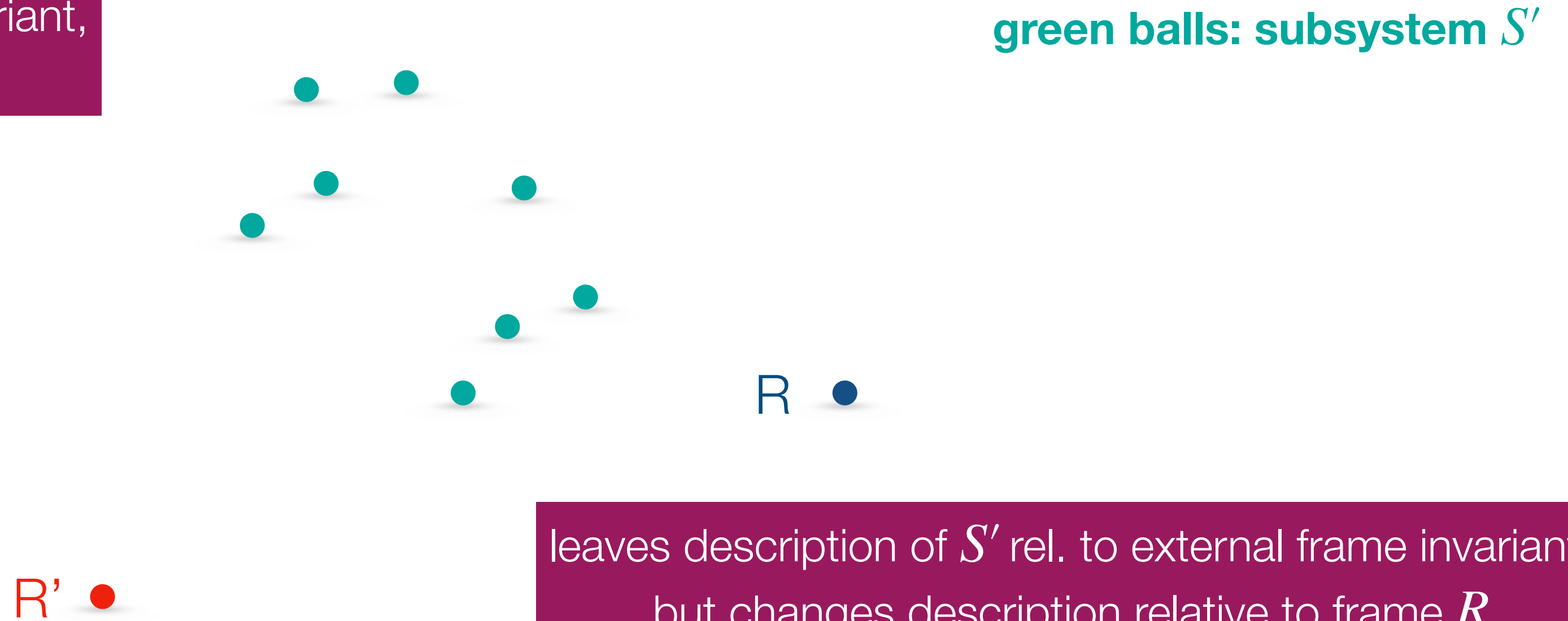
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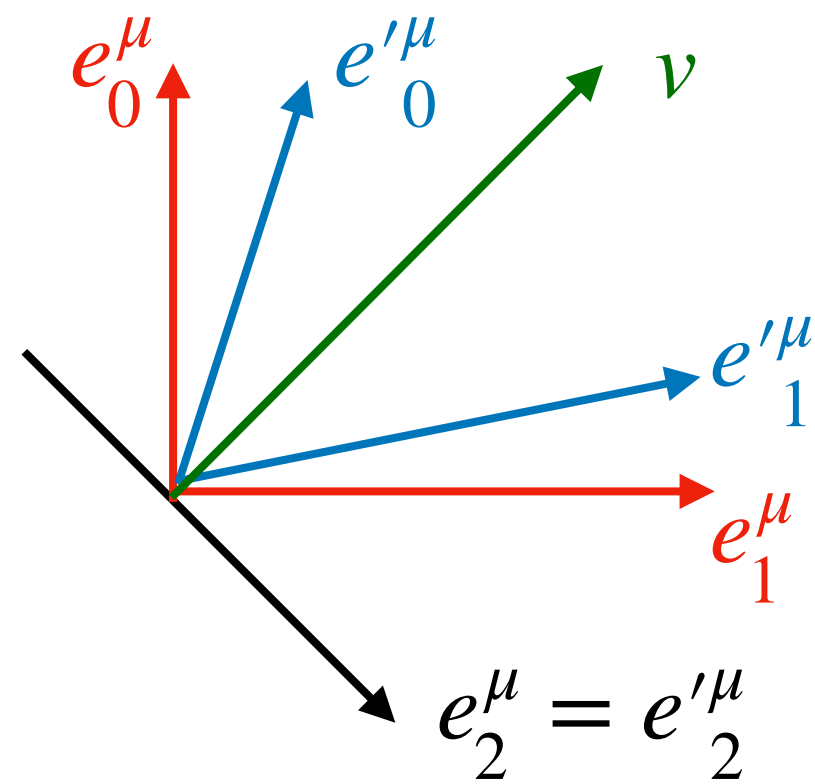
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1. kinematical and relational (gauge inv.) notion of subsystem **distinct**
2. gauge inv. notion of subsystem depends on choice of RF
 \Rightarrow gauge inv. correlations, thermal properties, ... are RF dependent

Warmup: Special relativity with internal frames



introduce second internal frame

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$$v_a = v^\mu \eta_{\mu\nu} e_a^\nu = v^\mu e'_{\mu a'} e'^{\nu a'} e_a^\nu = v_{a'} \Lambda^{a'}_a$$

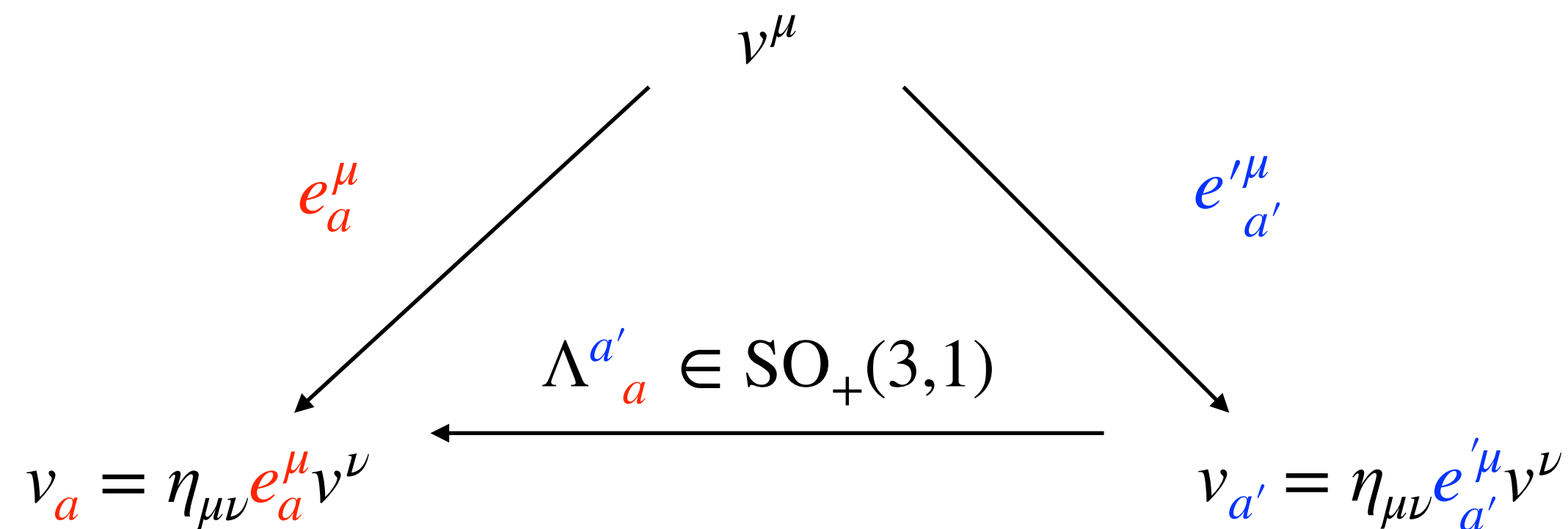
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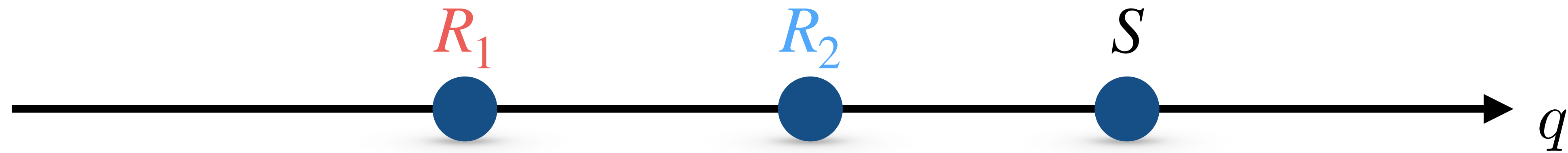


subsystem relativity \Rightarrow relativity of simultaneity

Quantum reference frames

... or frames in relative superposition

Example: spatial QRFs in 1D

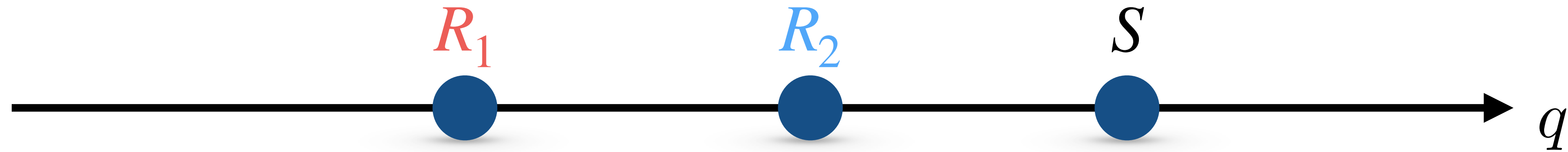


setup relative to external (possibly fictitious) frame:

$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

space of externally distinguishable states

Example: spatial QRFs in 1D



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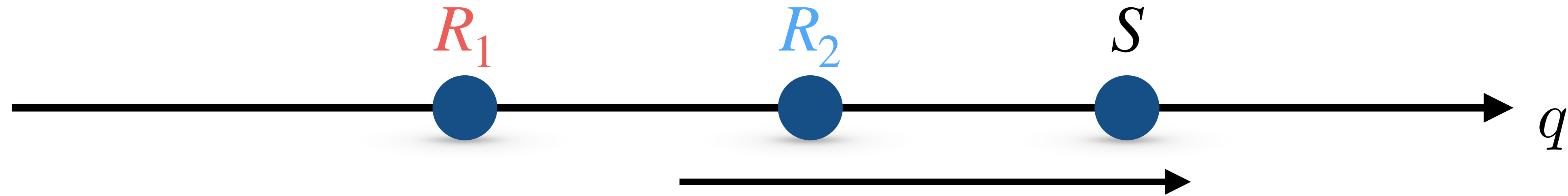
space of externally distinguishable states

global translations as external frame transformations, i.e. **gauge transformations**

$$U_{R_1 R_2 S}(x) = e^{ix(p_1 + p_2 + p_S)}$$

analog of $\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$ in SR

Example: spatial QRFs in 1D



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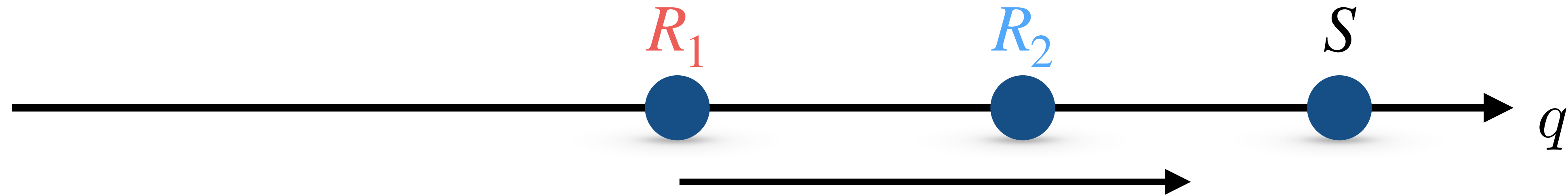
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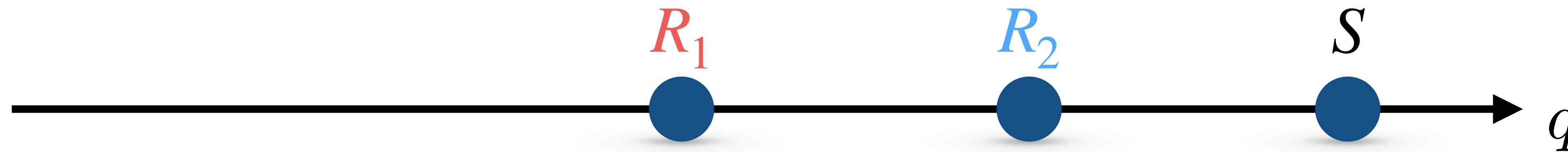
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frame orientation states for R_1 $|q\rangle_{R_1}$

analog of tetrad $e^\mu_a \in \text{SO}_+(3,1)$ in SR

frame reorientations (only act on frame)

$$U_{R_1}(x) = e^{ixp_1}$$

analog of $\Lambda^a_b e^\mu_a$ in SR

\Rightarrow 2 commuting group actions (here trivial)

Example: spatial QRFs in 1D



recall relational observables from SR $v_a = v^\mu \eta_{\mu\nu} e_a^\nu$

frame-orientation conditional gauge transformation

Relational observables through G -twirl:

$$O_{f_{R_2S}, R_1}(x) = \int dq U_{R_1R_2S}(q) \left(|x\rangle\langle x|_{R_1} \otimes f_{R_2S} \right) U_{R_1R_2S}^\dagger(q)$$

“what’s the value of f_{R_2S} when R_1 is in position x ?”

frame-orientation conditional gauge transformation (controlled unitary)

Example: spatial QRFs in 1D



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gauge-inv. $[O_{f_{R_2S}, R_1}(x), U_{R_1R_2S}(y)] = 0$

Example: spatial QRFs in 1D



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gauge-inv. $[O_{f_{R_2S}, R_1}(x), U_{R_1R_2S}(y)] = 0$

for example, get $O_{q_2, R_1}(0) = q_2 - q_1$ and $O_{q_S, R_1}(0) = q_S - q_1$

Perspective-neutral formulation of QRF covariance



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

space of externally distinguishable states

gauge-inv. (external frame-indep.) physical Hilbert space

$$\mathcal{H}_{\text{phys}} \quad \text{with} \quad |\psi\rangle_{\text{phys}} = U_{R_1 R_2 S}(x) |\psi\rangle_{\text{phys}}$$

space of internally distinguishable states

Perspective-neutral formulation of QRF covariance



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$$\mathcal{H}_{\text{phys}} \quad \text{with} \quad |\psi\rangle_{\text{phys}} = U_{R_1 R_2 S}(x) |\psi\rangle_{\text{phys}}$$

states such that

$$(p_1 + p_2 + p_S) |\psi\rangle_{\text{phys}} = 0$$

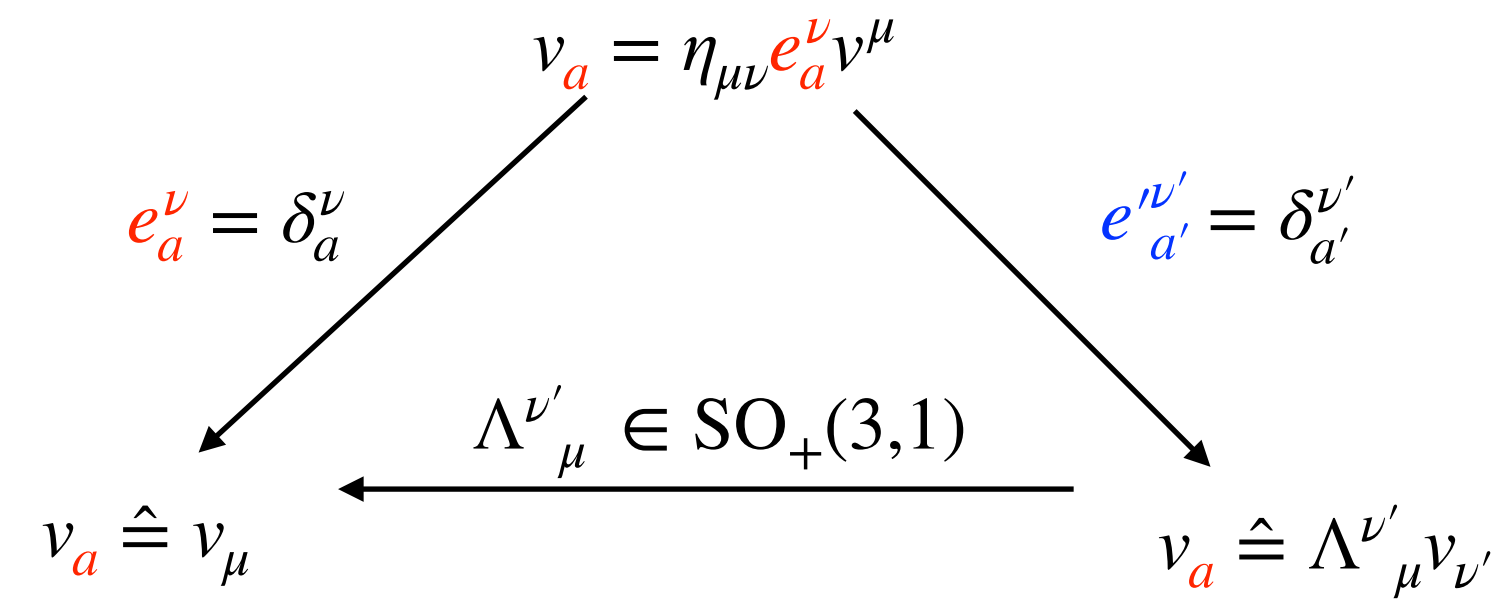
space of internally distinguishable states

⇒ redundancy: many different ways in describing same invariant $|\psi_{\text{phys}}\rangle$

⇒ associate with different internal QRF choices: redundant = reference DoFs

$\mathcal{H}_{\text{phys}}$ is a (internal frame) **perspective-neutral space**: description of physics prior to having chosen internal reference system relative to which remaining DoFs are described

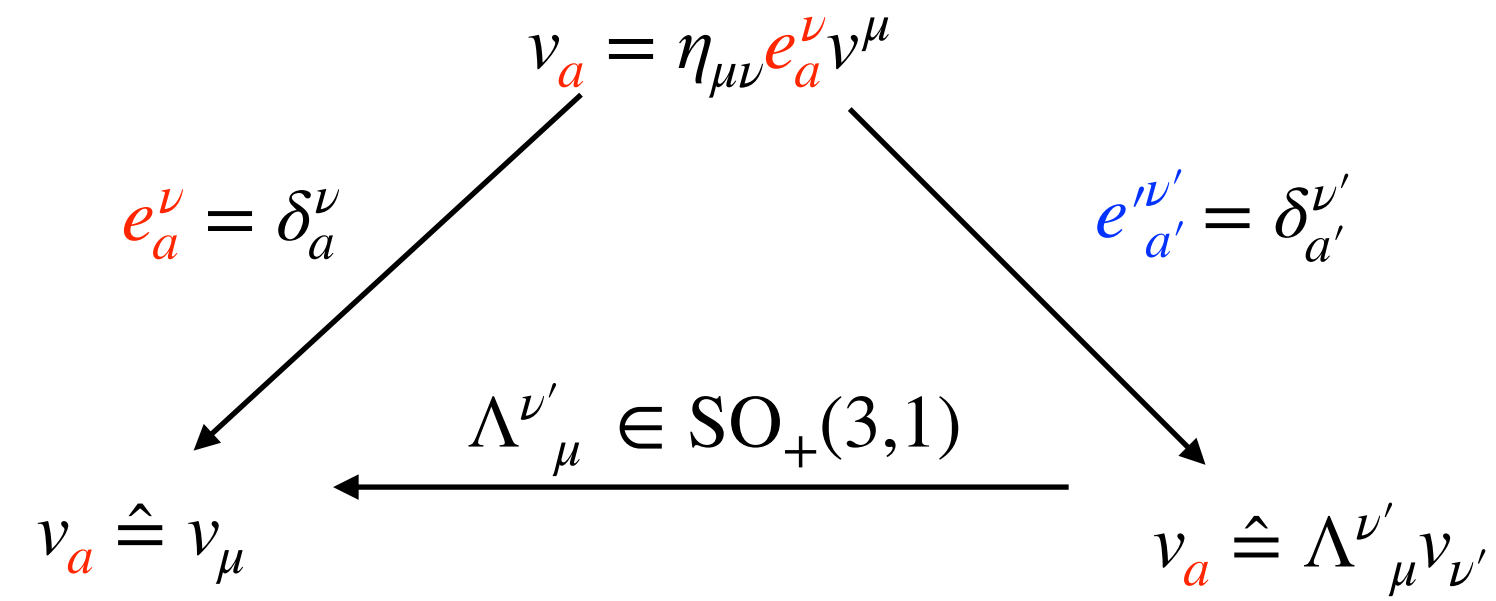
gauge-induced QRF changes: quantum coordinate changes



recall: “jumping into frame perspective” through gauge-fixing

gauge-induced QRF changes: quantum coordinate changes

Vanrietvelde, PH, Giacomini, Castro Ruiz, Quantum 2020



recall: "jumping into frame perspective" through gauge-fixing

perspective-neutral

$\mathcal{H}_{\text{phys}}$

$$\varphi_{R_1} = \langle q = 0 |_{R_1} \otimes I_{R_2 S}$$

$$\varphi_{R_2} = \langle q = 0 |_{R_2} \otimes I_{R_1 S}$$

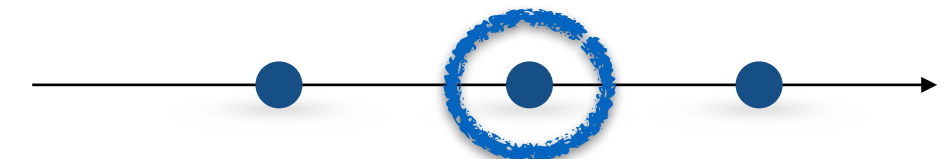
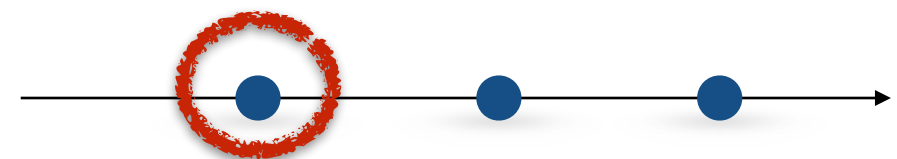
R_1 perspective

R_2 perspective

$$\mathcal{H}_{R_2 S | R_1} = L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

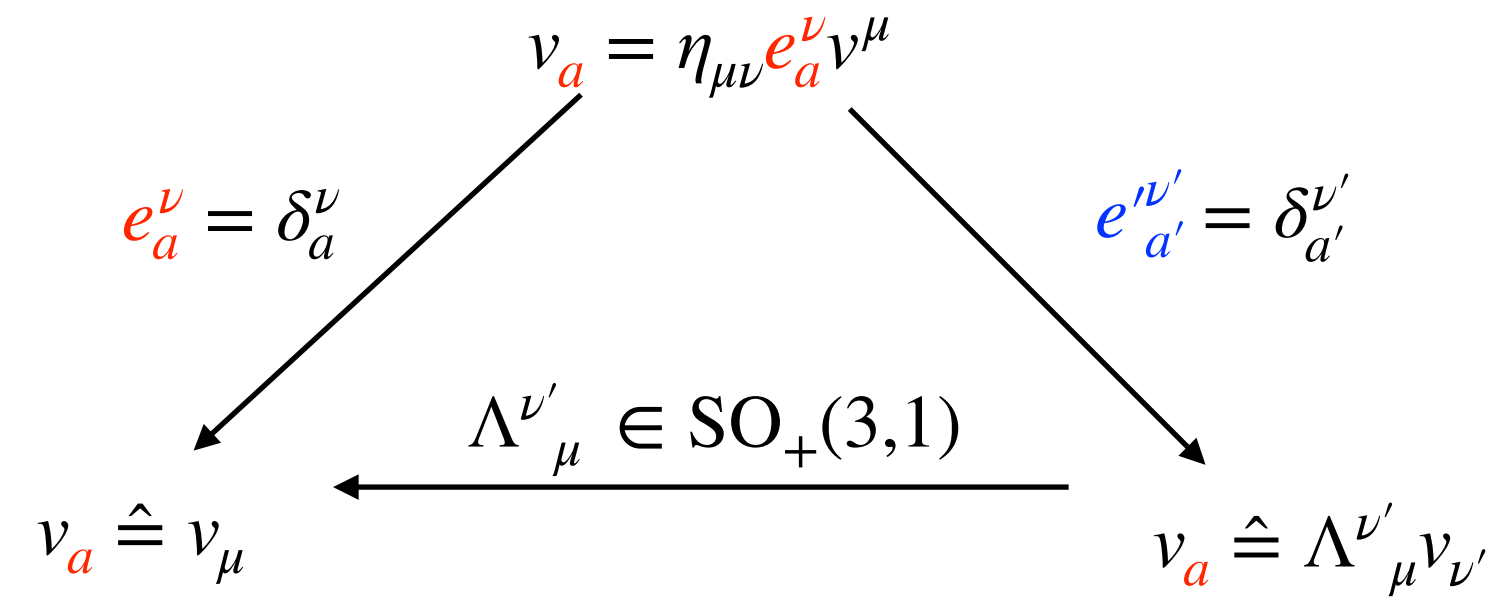
$$\mathcal{H}_{R_1 S | R_2} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_S$$

$$V_{R_1 \rightarrow R_2} = \varphi_{R_2} \circ \varphi_{R_1}^{-1}$$



gauge-induced QRF changes: quantum coordinate changes

Vanrietvelde, PH, Giacomini, Castro Ruiz, Quantum 2020



recall: "jumping into frame perspective" through gauge-fixing

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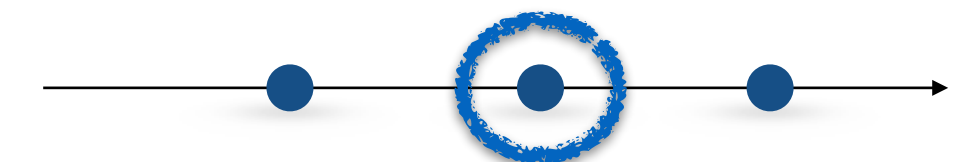
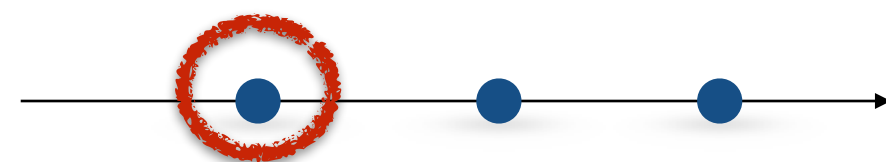
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R_1 perspective

R_2 perspective

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$$\mathcal{H}_{R_1 S | R_2} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_S$$



$$V_{R_1 \rightarrow R_2} = \varphi_{R_2} \circ \varphi_{R_1}^{-1}$$

$$= \mathbb{F}_{12} \int dq | -q \rangle \langle q |_{R_2} \otimes U_S(q)$$

reproduces QRF transf. from earlier

[Giacomini et al Nat. Comm. '19]

gauge-induced QRF changes: quantum coordinate changes

QRF relativity of entanglement and superposition

perspective-neutral

$\mathcal{H}_{\text{phys}}$

$$\varphi_{R_1} = \langle q = 0 |_{R_1} \otimes I_{R_2 S}$$

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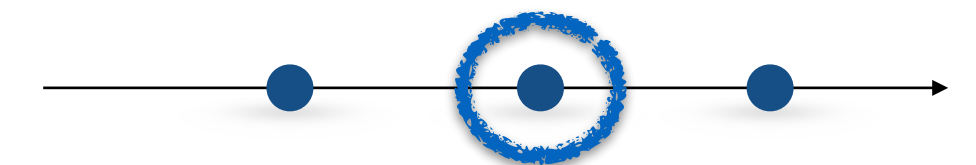
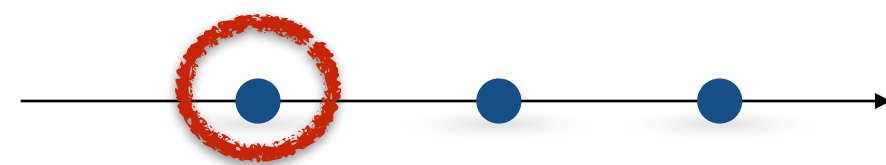
R_1 perspective

R_2 perspective

$$\mathcal{H}_{R_2 S | R_1} = L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

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$$V_{R_1 \rightarrow R_2} = \mathbb{F}_{12} \int dq | -q \rangle_{R_2} \langle q |_{R_1} \otimes U_S(q)$$



$$|q_1\rangle_{R_2} \otimes |x\rangle_S + |q_2\rangle_{R_2} \otimes |x\rangle_S$$

$$| -q_1 \rangle_{R_1} \otimes |x - q_1\rangle_S + | -q_2 \rangle_{R_1} \otimes |x - q_2\rangle_S$$

QRF relativity of subsystems

QRF perspectives φ_{R_1} and φ_{R_2} are nothing but two tensor product structures on gauge-inv. $\mathcal{H}_{\text{phys}}$

perspective-neutral

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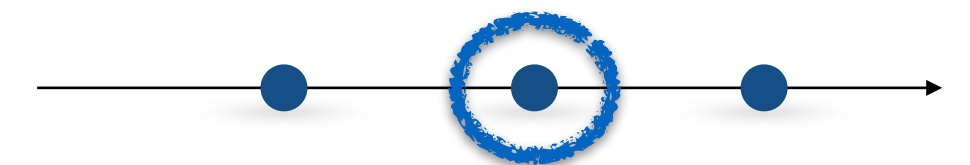
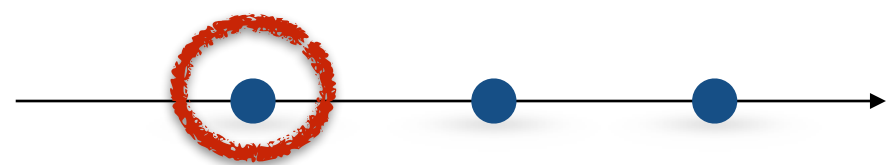
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\Rightarrow inequivalent because QRF transf. $V_{R_1 \rightarrow R_2}$ a nonlocal unitary

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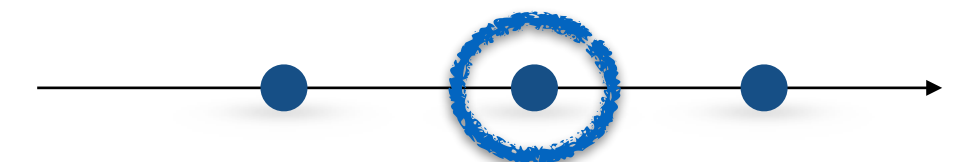
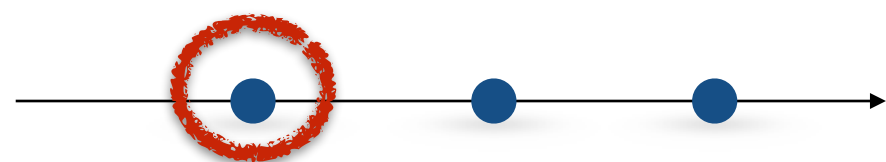
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QRF relativity of subsystems

different factorization of total gauge-inv. algebra into commuting subalgebras

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perspective-neutral

generated by canonical pairs $(q_S - q_1, p_S)$ and $(q_2 - q_1, p_1)$ which become local in

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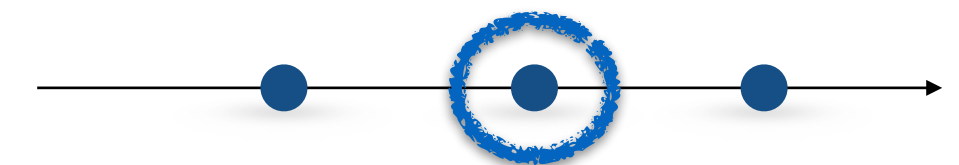
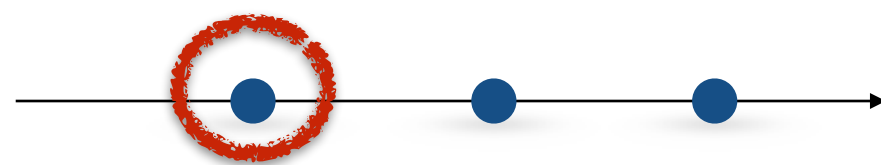
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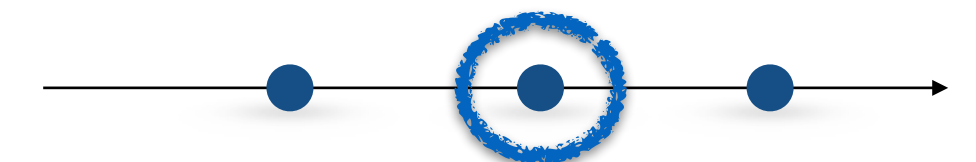
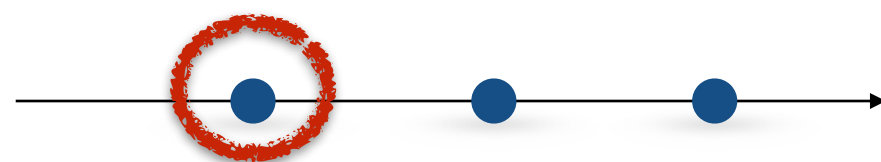
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have $\mathcal{A}_{S|R_2} \neq \mathcal{A}_{S|R_1}$

realization of virtual subsystems, Zanardi '00s

QRFs for general unimodular Lie groups

de la Hamette, Galley, PH, Loveridge, Müller
2110.13824;

works similarly, essentially

global translations $U_{R_1 R_2 S}(x)$ \longrightarrow **gauge transf.** $U_{R_1}(g) \otimes U_{R_2}(g) \otimes U_S(g)$

frame orientation states $|q\rangle_R$ \longrightarrow **coherent states** $|\phi(g)\rangle_R$

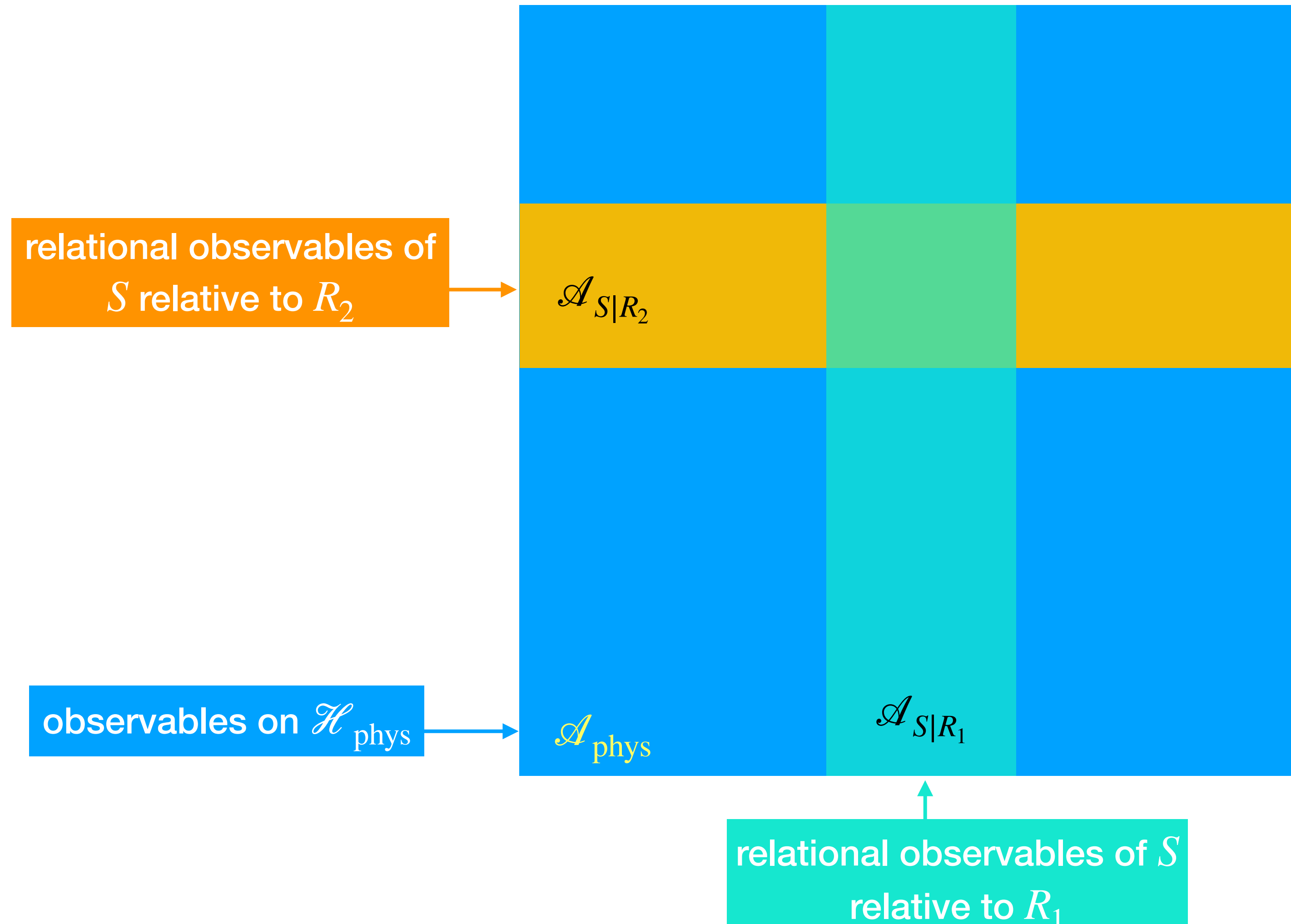
spatial integration $\int dq$ \longrightarrow **group integration** $\int_G dg$

.....

Quantum relativity of subsystems

Ahmad, Galley, PH, Lock, Smith PRL '22;
de la Hamette, Galley, PH, Loveridge, Müller
2110.13824;
Kotecha, Mele, PH to appear

3 kinematical subsystems $\mathcal{H}_{\text{kin}} = \mathcal{H}_{R_1} \otimes \mathcal{H}_{R_2} \otimes \mathcal{H}_S$



Recall: kinematical vs. relational subsystems

leaves description of S' relative to R invariant,
but changes it relative to R'

green balls: subsystem S'



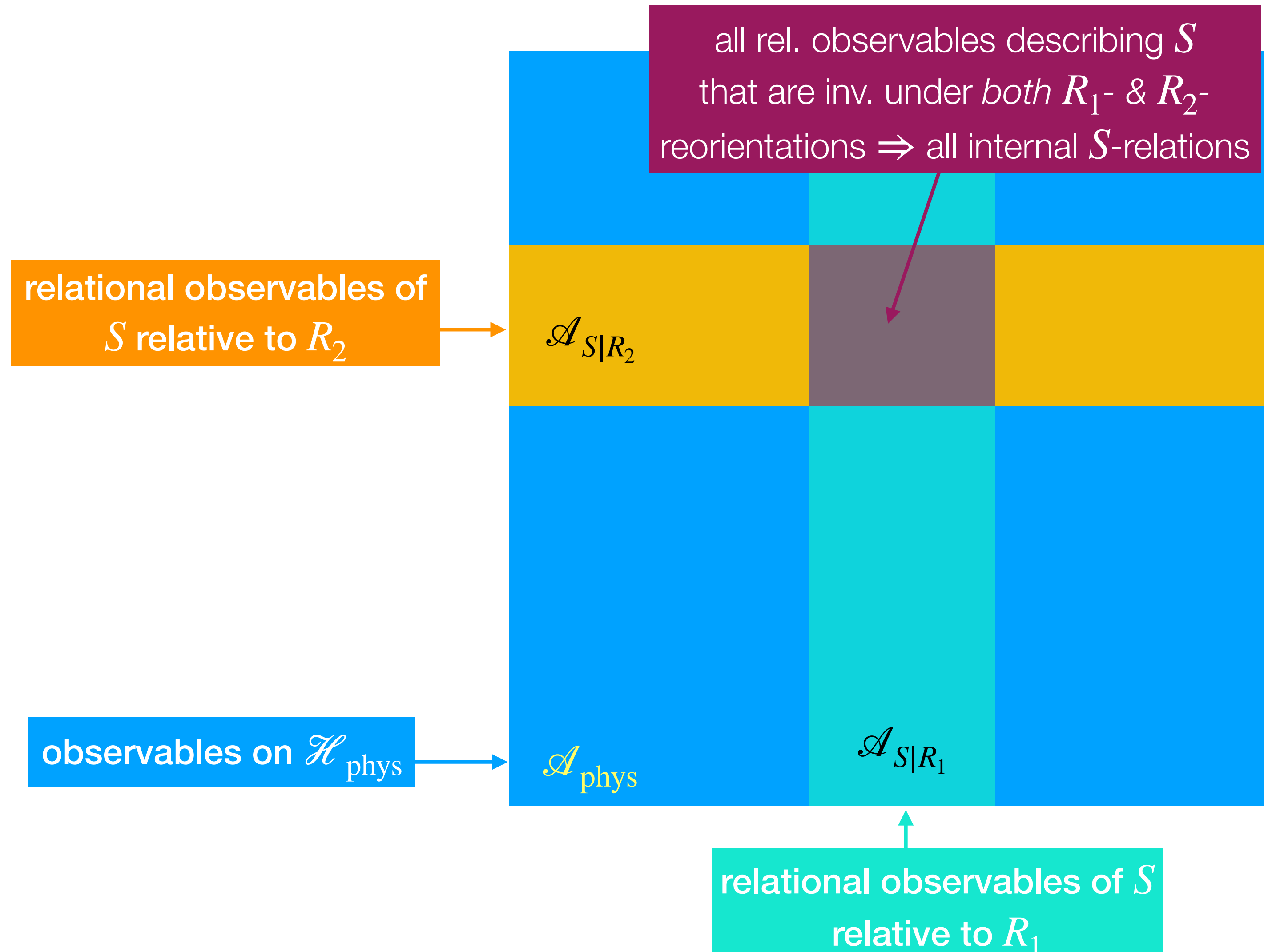
leaves description of S' rel. to external frame invariant,
but changes description relative to frame R

\Rightarrow overlap of rel. observable algebras $\mathcal{A}_{S|R} \cap \mathcal{A}_{S|R'} = \{\text{internal rel. obs. of } S\}$ (but don't coincide)

Quantum relativity of subsystems

Ahmad, Galley, PH, Lock, Smith PRL '22;
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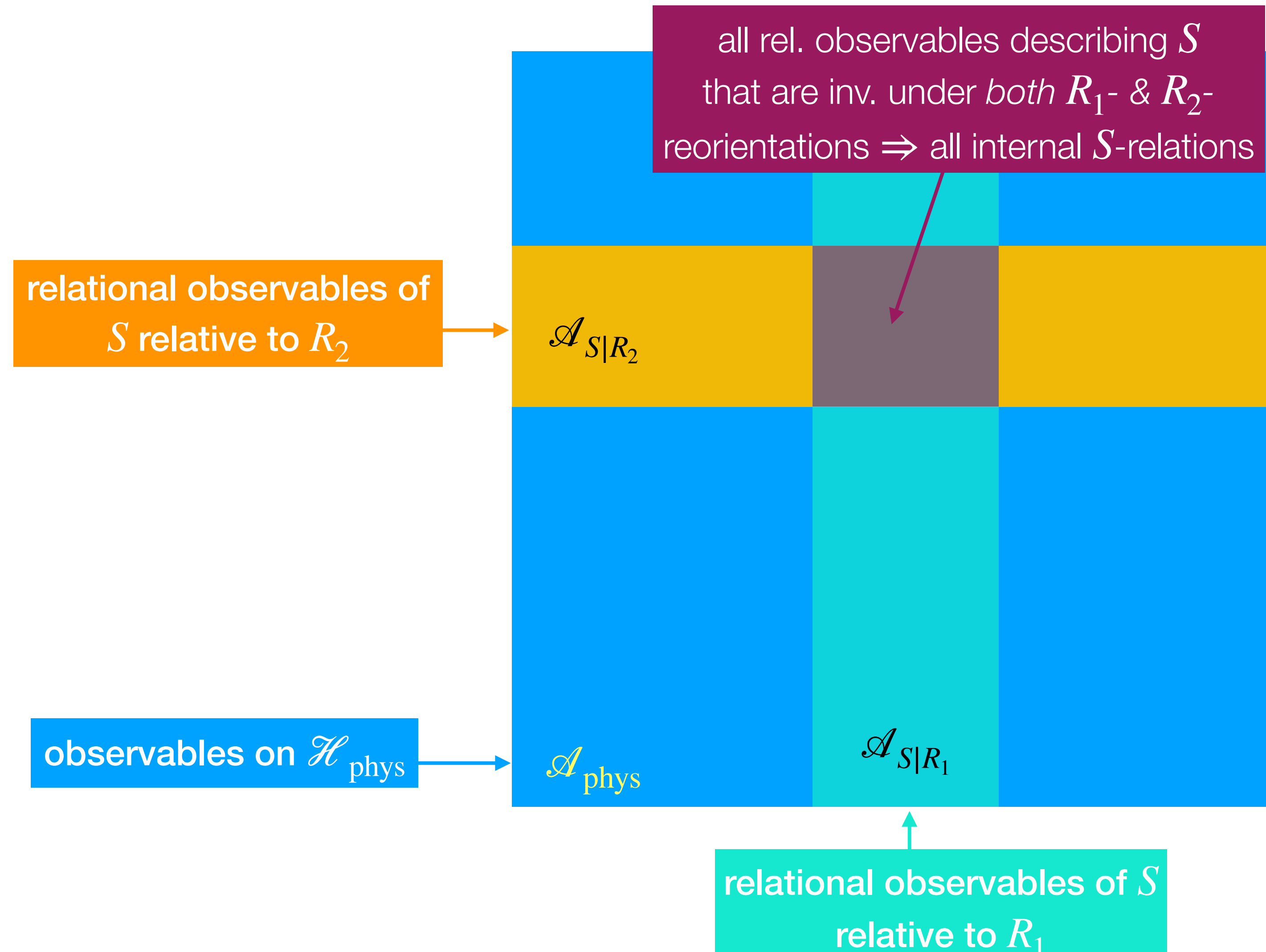
Ahmad, Galley, PH, Lock, Smith PRL '22;
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2110.13824;
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3 kinematical subsystems $\mathcal{H}_{\text{kin}} = \mathcal{H}_{R_1} \otimes \mathcal{H}_{R_2} \otimes \mathcal{H}_S$

⇒ different relational observable subalgebras
inside total invariant algebra

⇒ induce different gauge-inv. tensor factorizations
(not in general factorization across $R_j | R_i$ and $S | R_i$)

⇒ different appearance of same physics



Upshot: frame-dependent gauge-inv. properties

“frames R_1 and R_2 mean different inv. DoFs when they refer to subsystem S ”

Ahmad, Galley, PH, Lock, Smith PRL '22;
de la Hamette, Galley, PH, Loveridge, Müller
2110.13824

if factorizability in two frame perspectives, i.e.

$$\mathcal{A}_{\text{phys}} \simeq \mathcal{A}_{S|R_1} \otimes \mathcal{A}_{R_2|R_1} \simeq \mathcal{A}_{S|R_2} \otimes \mathcal{A}_{R_1|R_2} \quad \text{but} \quad \mathcal{A}_{S|R_2} \neq \mathcal{A}_{S|R_1}$$

then correlations/entanglement of S with its complement will in general differ in two perspectives

(see also Giacomini, Castro-Ruiz, Brukner '19; Castro-Ruiz, Oreshkov '21)

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then correlations/entanglement of S with its complement will in general **differ** in two perspectives

(see also Giacomini, Castro-Ruiz, Brukner '19; Castro-Ruiz, Oreshkov '21)

⇒ gauge-inv. entanglement entropy in general $S(\rho_{S|R_2}) \neq S(\rho_{S|R_1})$ for **same global physical state**

⇒ dynamics of S can be closed/isolated relative to R_1 and open relative to R_2

(can map unitary dynamics/zero heat & work exchange into
open dynamics/ non-zero heat & work exchange in other perspective)

Kotecha, Mele, PH to appear

⇒ **QRF relativity of superpositions, correlations, equilibrium, thermodynamics, ...**

Conclusions

- **Natural extension of relativity principle into quantum realm**
based on internal QRFs \Rightarrow in terms of group structures really the same as in SR
- **Systematic method for changing QRF perspectives**
accommodates RFs in relative superposition
- **Gauge-inv. subsystems depend on choice of QRF (“quantum relativity of subsystems”)**
 \Rightarrow correlations, thermal properties, dynamics, depend on frame

\Rightarrow works completely analogously with (so far classical) dynamical frames in gauge theory and gravity

Appendix

Symmetry-induced QRF changes

de la Hamette, Galley, PH, Loveridge, Müller 2110.13824

changes of relational observables, recall:

$$v_a = v_{a'} \Lambda_{a'}^{a'}$$

relational observable rel. to e

relational observable rel. to e'

- RF transformation between two frames is $\Lambda_{a'}^{a'} = e_{\mu}^{a'} e_a^{\mu} \in \text{SO}_+(3,1)$
relational observable describing 1st rel. to 2nd frame

relation-conditional frame reorientation

Symmetry-induced QRF changes

de la Hamette, Galley, PH, Loveridge, Müller 2110.13824

can do analog in QT: G-twirl for symmetries $V_{R_1}(g) \otimes \mathbf{1}_{R_2S}$

relation-conditional frame reorientation

